

DEFORMATIONAL AND PERTURBATIVE RIGIDITY NEAR DE LA LLAVE EXAMPLES

ANDREY GOGOLEV AND MARTIN LEGUIL

ABSTRACT. De la Llave's examples are Anosov diffeomorphisms on the four-torus $\mathbb{T}^4$ with constant Lyapunov spectrum, yet they are not C^1 -conjugate to the linear model or to each other. Nevertheless, we show that such examples are “locally exceptional”: we prove deformation and local rigidity for generic diffeomorphisms in proximity of de la Llave's examples.

1. INTRODUCTION

This paper is devoted to deformational and perturbative rigidity of Anosov diffeomorphisms in dimension four. Consider an Anosov diffeomorphism F_0 on a compact Riemannian manifold M . In other words, there exists a DF_0 -invariant splitting $TM = E^s \oplus E^u$ of the tangent bundle into stable and unstable spaces E^s , E^u , such that vectors in E^s (resp. E^u) get exponentially contracted under forward (resp. backward) iteration of DF_0 . By structural stability, for any diffeomorphism G which is sufficiently C^1 -close to F_0 , there exists a homeomorphism h , called a (topological) *conjugacy*, such that $h \circ F_0 = G \circ h$. Yet, there exist plenty of obstructions to the existence of a C^1 conjugacy; indeed, if the conjugacy map h above can be chosen to be C^1 , then for any periodic point $p = F_0^n(p)$, we can differentiate the conjugacy equation and obtain

$$Dh(p)DF_0^n(p)(Dh(p))^{-1} = DG^n(h(p)),$$

that is, the differentials $DF_0^n(p)$ and $DG^n(h(p))$ must be conjugate. In fact, for any periodic point p , it is easy to arrange the perturbation G such that the above condition fails (and this is typical). Therefore, there are countably many obstructions to lift to hope for the existence of a C^1 conjugacy, which are associated to the periodic points.

We say that F_0 and G are *isospectral* if for every periodic point $p = F_0^n(p)$ the linearized return maps $DF_0^n(p)$ and $DG^n(h(p))$ have the same collection of eigenvalues. In particular, given a one-parameter family¹ of Anosov diffeomorphisms $\{F_s\}_{s \in [0,1]}$, there exists a unique continuous family of conjugacies $\{h_s\}_{s \in [0,1]}$ such that $h_s \circ F_0 = F_s \circ h_s$ and $h_0 = \text{id}$. Accordingly, such family is called *isospectral* if for every periodic point $p = F_0^n(p)$, all the linearized return maps $DF_s^n(h_s(p))$, $s \in [0,1]$, have the same collection of eigenvalues.

It is a classical result that for 2-dimensional Anosov diffeomorphisms, the collection of eigenvalues of periodic points is actually a complete set of moduli of smooth conjugacy classes:

Theorem 1.1 (Marco-Moriyón, de la Llave [MM87a, DIL87, MM87b, DIL92]). *Let F_0, G be two C^r , $r \in (1, \infty] \cup \{\omega\}$, Anosov diffeomorphisms on $\mathbb{T}^2$ which are topologically conjugated, and whose eigenvalues at corresponding periodic points match. Then the conjugacy is C^{r*} regular, with*

$$r_* = r, \text{ if } r \notin \mathbb{N}, \text{ and } r_* = (r - 1) + \text{Lip}, \text{ if } r \in \mathbb{N}.$$

In particular, in this low-dimensional situation, once the conjugacy is C^1 , due to the bootstrap phenomenon, it is actually as regular as the Anosov diffeomorphisms themselves. In

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¹We assume continuity in C^2 topology.

higher dimension, it is not always possible to bootstrap the regularity of a C^1 conjugacy between Anosov diffeomorphisms. The lack of bootstrap phenomenon is typically related to the lack of regularity of the leaves of intermediate (weak stable/unstable) foliations, when they exist.

Definition 1.2. A C^r Anosov diffeomorphism F_0 is called C^r -*locally rigid*, $r \geq 2$, if there exists a neighborhood $\mathcal{U}$ of F_0 in the C^r topology such that for any diffeomorphism $G \in \mathcal{U}$ that is isospectral to F_0 , the corresponding conjugacy h is, in fact, a C^{1+H} diffeomorphism.² Similarly, F_0 is called C^r -*deformation rigid* if for any isospectral one-parameter family $\{F_s\}_{s \in [0,1]}$ based at F_0 such that the map $s \mapsto F_s$ is continuous in the C^r topology, the corresponding family of conjugacies h_s is, in fact, a family of C^{1+H} diffeomorphisms.

In higher dimension, there exist counterexamples to the deformation and local rigidity of Anosov diffeomorphisms constructed by de la Llave [DIL92], which we now proceed to recall.

1.1. de la Llave examples. Let A and B be automorphisms of the 2-torus $\mathbb{T}^2$ induced by hyperbolic matrices in $\mathrm{SL}(2, \mathbb{Z})$. We will assume that the smaller eigenvalues λ and μ of A and B , respectively, satisfy the following inequalities: $0 < \lambda < \mu < 1$. Define $\alpha = \log \mu / \log \lambda$ and notice that $\alpha \in (0, 1)$.

Definition 1.3. A *de la Llave diffeomorphism* $L_{\varphi_0}: \mathbb{T}^4 \rightarrow \mathbb{T}^4$ is defined as a skew-product

$$L_{\varphi_0}(x, y) = (Ax, By + \varphi_0(x)), \quad \forall (x, y) \in \mathbb{T}^2 \times \mathbb{T}^2,$$

where $\varphi_0: \mathbb{T}^2 \rightarrow \mathbb{T}^2$ is a smooth function.

Such diffeomorphism is Anosov and, if φ_0 is homotopic to a constant, is conjugate to the linear product automorphism L_0 . More generally, if φ_1 is homotopic to φ_0 then the corresponding de la Llave diffeomorphism L_{φ_1} is conjugate to L_{φ_0} with conjugacy h given by

$$(1.1) \quad h(x, y) = (x, y + \psi(x)), \quad \forall (x, y) \in \mathbb{T}^2 \times \mathbb{T}^2,$$

where $\psi: \mathbb{T}^2 \rightarrow \mathbb{T}^2$ is null-homotopic. Indeed, a solution of the form in (1.1) to the conjugacy equation $h \circ L_{\varphi_0} = L_{\varphi_1} \circ h$ can be obtained explicitly by solving the cohomological equation for $\varphi_1 - \varphi_0$:

$$(1.2) \quad (\varphi_1 - \varphi_0)(x) = \psi \circ A(x) - B \circ \psi(x), \quad \forall x \in \mathbb{T}^2.$$

We can check that this equation has a unique solution $\psi \in C^\alpha(\mathbb{T}^2)$, and then, the associated map h is the unique conjugacy between L_{φ_0} and L_{φ_1} in the homotopy class of the identity, and it is also C^α regular. However, for some simple choices of $\varphi_1 - \varphi_0$, and even for a generic choice, the function ψ is not $C^{\alpha+\varepsilon}$ regular, for any $\varepsilon > 0$. Let us, for example, consider the case where $\varphi_1 - \varphi_0$ is a single Fourier mode $x \mapsto e_u \cos(2\pi \langle k_0, x \rangle)$, $k_0 \in \mathbb{Z}^2 \setminus \{(0, 0)\}$, where $\langle \cdot, \cdot \rangle$ is the function on $\mathbb{T}^2 \times \mathbb{T}^2$ induced by the Euclidean inner product on $\mathbb{R}^2$, and e_u is an unstable eigenvector for B , $Be_u = \mu^{-1}e_u$. The solution ψ to (1.2) is of the form $\tilde{\psi}e_u$, for some function $\tilde{\psi}: \mathbb{T}^2 \rightarrow \mathbb{R}$. Considering the Fourier series $\sum_{k \in \mathbb{Z}^2} \hat{\psi}_k e^{i2\pi \langle k, x \rangle}$ of $\tilde{\psi}$, we find that

$$\hat{\psi}_{\pm(A^\top)^n k_0} = -\mu^{n+1}, \quad \forall n \geq 0, \quad \hat{\psi}_k = 0, \quad \forall k \notin \{\pm(A^\top)^n k_0\}_{n \geq 1} \cup \{(0, 0)\}.$$

For some constant $K > 0$, we have $|(A^\top)^n k_0| \geq K\lambda^{-n}$, hence the decay of the Fourier coefficients ensures that $\tilde{\psi} \in C^\alpha$ but $\tilde{\psi} \notin C^{\alpha+\varepsilon}$, for any $\varepsilon > 0$. Consequently, the conjugacy h is merely C^α regular despite the fact that all eigenvalues at corresponding periodic points (and Lyapunov exponents of all corresponding measures) of L_{φ_0} and L_{φ_1} are the same. This very interesting phenomenon was discovered and explained by de la Llave in [DIL92], see also [Gog08, GRH23] for later expositions. De la Llave diffeomorphisms also demonstrate several other interesting dynamical features. For example, they are non-linear Anosov diffeomorphisms for which the invariant volume measure has maximal entropy; they could be used to exhibit insufficient regularity of weak foliations (when $\alpha > 1$), etc.

²Here, C^{1+H} means C^1 with Hölder continuous differential.

Furthermore, by linearly extrapolating into a family of diffeomorphisms $\{L_{\varphi_s}\}_{s \in [0,1]}$, $\varphi_s = \varphi_0 + s(\varphi_1 - \varphi_0)$, $s \in [0,1]$, we obtain an isospectral family, yet the conjugacy between any two diffeomorphisms in the family is merely C^α .

The first author studied local C^{1+H} -conjugacy classes of de la Llave diffeomorphisms in [Gog08]. In particular, it was established that a C^{1+H} isospectral perturbation of a de la Llave diffeomorphism L_{φ_0} is C^{1+H} -conjugate to a de la Llave diffeomorphism L_{φ_1} , where φ_1 is C^1 -close to φ_0 . Hence, de la Llave examples completely absorb the failure of periodic eigenvalue rigidity for the product automorphism L_0 . This naturally raises the question of description of local C^{1+H} -conjugacy classes for perturbations of L_0 and, more generally, of L_{φ_0} .

1.2. Deformation and perturbative rigidity. The following theorems constitute progress on a smooth classification in proximity of de la Llave diffeomorphisms. Recently, Rafael de la Llave communicated to the authors that he and others considered the possibility of such rigidity in proximity of the examples at the time when examples were discovered.

Theorem A. Let L_φ be a de la Llave diffeomorphism. Then there exist a C^1 small neighborhood $\mathcal{U}$ of L_φ and a C^2 -open C^∞ -dense subset $\mathcal{V} \subset \mathcal{U}$ such that each Anosov diffeomorphism $F \in \mathcal{V}$ is C^2 -deformation rigid.

Theorem B. Let L_φ be a de la Llave diffeomorphism. Then there exist a C^1 small neighborhood $\mathcal{U}$ of L_φ and a C^2 -open C^∞ -dense subset $\mathcal{V} \subset \mathcal{U}$ such that each Anosov diffeomorphism $F \in \mathcal{V}$ is C^2 -locally rigid.

Remark 1.4. We note that there is no obvious way to deduce Theorem A from Theorem B. The issue is that an isospectral deformation is not required to be local and stay in a neighborhood of F (moreover, can leave the neighborhood of L_φ) and, hence, one cannot apply Theorem B to the path, but only to the beginning of the path. One then could attempt covering the path by a finite number of small neighborhoods and gradually conjugate the whole path to F . However, this approach is problematic since the conjugacy we obtain is only C^{1+H} regular, and it seems hard to make our proof work in this low regularity predicament.

Remark 1.5. To keep the exposition clean, we work with C^∞ diffeomorphisms. But by carefully inspecting the proof, one can check that the same results are true for C^4 diffeomorphisms.

The latter result suggests the following open questions.

- (Q1) *Global rigidity.* Show that for $F \in \mathcal{V}$, any Anosov diffeomorphism G which is isospectral to F is C^1 -conjugate to F .
- (Q2) *Volume-preserving version.* Show that there exists a C^1 small neighborhood $\mathcal{U}^{\text{vol}}$ of L_φ in the space of volume-preserving Anosov diffeomorphisms and a C^2 -open C^∞ -dense subset $\mathcal{V}^{\text{vol}} \subset \mathcal{U}^{\text{vol}}$ such that each $F \in \mathcal{V}^{\text{vol}}$ is (C^2 -locally) rigid.
- (Q3) *Jacobian rigidity.* Show that one can weaken the isospectrality assumption in Theorems A and B to mere matching of Jacobians at corresponding periodic points.

Remark 1.6. Let us make some observations on these questions:

- our strategy relies on certain period expansions near periodic points with dissipative behavior along the center (see e.g. Proposition 3.1); in the volume-preserving case, as suggested in Question (Q2), such expansions look different (see [GLH25, Proposition 4.17]), and the leading exponential term mixes data coming from the (weak-)stable and unstable directions, which would make harder the identification of bifurcations such as those we research, e.g., in Proposition 3.15; also, considering suspensions via logarithmic full Jacobian is not useful in the volume-preserving setting;
- regarding Question (Q3), while the expansions we consider here do assume matching of eigenvalues along the different directions, we may be able to derive such information from matching of the full Jacobian, similarly to what was achieved in [GLH25, Theorem A-Corollary B].

In fact, some years ago the first author gave a conjectural description of all possible “generalized de la Llave” examples. While our results give a definitive progress on rigidity near de la Llave examples, they still fall short of such an explicit description. We record this conjecture here.

Conjecture 1.7. *Let L_φ be a de la Llave diffeomorphism and let $\mathcal{U}$ be a small C^1 -neighborhood of it where the 4-way dominated splitting survives. Let $F \in \mathcal{U}$ and assume that the center foliation W_F^c by 2-tori, which is the integral foliation of the weak stable and weak unstable distributions of F , is not C^1 regular as a foliation.³ Then F is C^1 -locally rigid.*

Remark 1.8. The condition “ W_F^c is not C^1 regular” can be fully understood in terms of periodic data of F . Namely, W_F^c is C^1 regular if and only if for any periodic point p and any other periodic point $q \in W_F^c(p)$ the strong stable and strong unstable Lyapunov exponents of p and q are equal. This fact can be derived via a similar argument to the characterization of C^1 weak unstable foliations on $\mathbb{T}^3$ given in [Gog12].

Let us briefly summarize the main steps of the proof of Theorem B. Fix a diffeomorphism $F \in \mathcal{V}$, and a diffeomorphism G which is sufficiently C^2 -close to F and isospectral to F .

- (1) We consider the suspension flows X^t and Y^t over F and G , respectively, with roof functions given by the logarithmic full Jacobian (+const); by the isospectrality condition, the resulting flows are conjugate, i.e., for some homeomorphism H , we have $H \circ X^t = Y^t \circ H$, for all $t \in \mathbb{R}$.
- (2) We fix two periodic points $p = X^T(p)$, $\tilde{p} = X^{\tilde{T}}(\tilde{p})$, such that $DX^T(p)$, $DX^{\tilde{T}}(\tilde{p})$ expands, resp. contracts volume along the “center”; given homoclinic points q and $\tilde{q}$ to p and $\tilde{p}$, respectively, we consider sequences (p_n) , $(\tilde{p}_n)$ of periodic points in the associated horseshoe and study the asymptotics of their periods as $n \rightarrow +\infty$ (see Proposition 3.1).
- (3) For suitable choice of $q, \tilde{q}$, we show that matching of the periods of p_n and $H(p_n)$, resp. $\tilde{p}_n$ and $H(\tilde{p}_n)$ for the flows X^t, Y^t forces the conjugacy H to send the strong stable manifold of p to the strong stable manifold of $H(p)$, resp. the strong unstable manifold of $\tilde{p}$ to the strong unstable manifold of $H(\tilde{p})$ (see Propositions 2.4-2.6). More specifically, the asymptotic formulae for the periods allow us to recover the coefficients (one coefficient for each choice of $q, \tilde{q}$) by the leading exponential terms. This countable collection of coefficients “have full knowledge” of (completely determine) the position of the strong unstable manifold of p (resp. $\tilde{p}$) inside the 2-dimensional unstable manifold. And matching of the coefficients forces matching of strong stable manifolds under H . This step constitutes the main novel technique of this paper.
- (4) We show that the above condition implies that H preserves strong stable and strong unstable foliations.
- (5) From the preservation of invariant foliations and the isospectrality condition, we conclude that the conjugacy H is C^{1+H} regular, and similarly for the conjugacy between the diffeomorphisms F and G .

In Appendix A, we present a rigidity theorem at the level of 5-dimensional Anosov flows to elucidate key mechanisms and offer additional insight into the proofs of Theorems A and B. Indeed, as outlined in the proof sketch for Theorems A-B, our approach centers on analyzing the conjugacy of 5-dimensional Anosov flows that are suspensions over diffeomorphisms near de la Llave’s examples. Some properties demonstrated in this specific context actually extend to a broader class of 5-dimensional Anosov flows; the purpose of Theorem C is to formalize this generalization. However, the proofs of Theorems A and B remain formally independent of Theorem C.

³See also Remark 1.9 for additional insights about why anomalous regularity of invariant foliations can be an issue in the study of rigidity questions.

Remark 1.9. Complementing Remark 1.8, the regularity of invariant foliations also appeared as a possible issue for rigidity in the context of 3-dimensional Anosov flows, see [GLH25, Theorem E]. In fact, it is directly related to one of the “genericity” conditions we impose in the definition of the set $\mathcal{V}$ in Theorems A-B. Indeed, we require the non-vanishing of a quantity related to certain “templates” of the suspension flows X^t and Y^t , which can be thought of as the temporal coordinate of some invariant distributions (see Section 2 for the precise definition of these objects, and Proposition 3.10 for the associated genericity condition). By following the approach developed in [GLH25], we claim that for those $F \in \mathcal{V}$ which are non-rigid, the aforementioned templates of the suspension flow X^t exhibit anomalous regularity, which itself reflects anomalous regularity of the associated distributions, as described in items (1)-(2) of Theorem C. In particular, the genericity condition we need to impose on $F \in \mathcal{V}$ is related to their Jacobian, to avoid falling into cases (1)-(2) in Theorem C.

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2. PREPARATIONS AND STRUCTURE OF THE PROOF

Let us fix an automorphism $L_0 = (A, B)$, and a de la Llave diffeomorphism L_φ as in Subsection 1.1. We first explain how the neighborhood $\mathcal{U}$ is chosen. The automorphism L_0 admits constant dominated splitting $TT^4 = E_0^{ss} \oplus E_0^{ws} \oplus E_0^{wu} \oplus E_0^{uu}$ according to the eigendirections corresponding to λ, μ, μ^{-1} and λ^{-1} , respectively. It is a standard exercise on the cone technique to check that a de la Llave diffeomorphism L_φ also admits a dominated splitting $E_\varphi^{ss} \oplus E_\varphi^{ws} \oplus E_\varphi^{wu} \oplus E_\varphi^{uu}$ with the same exponential rates as L_0 . In fact $E_\varphi^{ws} = E_0^{ws}$ and $E_\varphi^{wu} = E_0^{wu}$, but the strong subbundles become different. The domination condition is C^1 -open, so we can choose an open set $\mathcal{U} \ni L_\varphi$ such that each $F \in \mathcal{U}$ is Anosov and admits a dominated splitting $E_F^{ss} \oplus E_F^{ws} \oplus E_F^{wu} \oplus E_F^{uu}$ which is partially hyperbolic, with 2-dimensional center $E_F^{ws} \oplus E_F^{wu}$. Let $E_F^s := E_F^{ss} \oplus E_F^{ws}$ and $E_F^u := E_F^{uu} \oplus E_F^{wu}$ be the full (2-dimensional) stable and unstable bundles. Since the decomposition $E_F^s \oplus E_F^u$ is hyperbolic, E_F^s and E_F^u integrate uniquely to stable and unstable foliations W_F^s and W_F^u , respectively. From partially hyperbolic theory (see [HPS77]) we have that the strong distributions E_F^{ss} and E_F^{uu} uniquely integrate to strong foliations W_F^{ss} and W_F^{uu} , respectively. Furthermore, a standard argument on the universal cover ensures that the weak distributions E_F^{ws} and E_F^{wu} integrate uniquely to weak foliations W_F^{ws} and W_F^{wu} , respectively [GG08]. Moreover, E_F^{ws} and E_F^{wu} integrate jointly to a “center” foliation by tori, which we denote by W_F^c , whose leaves are subfoliated by the leaves of W_F^{ws} and W_F^{wu} ; this fact can be established by comparison of divergence rates on the universal cover.⁴

We will need the above foliations to persist through an isospectral deformation which could, *a priori*, leave the neighborhood $\mathcal{U}$. First, by choosing $\mathcal{U}$ even smaller if needed, we can guarantee that $F \in \mathcal{U}$ has narrow periodic data (that is, the Lyapunov exponents at periodic points of F are sufficiently close to the Lyapunov exponents of L_φ). If a diffeomorphism G is conjugate to F and has the same periodic data (such as a diffeomorphism from an isospectral deformation of F) then G also admits a dominated splitting by the main result of [DG24].

⁴In fact, if φ is sufficiently C^1 -small, we can also appeal to the general Hirsch-Pugh-Shub machinery: since the center distribution $E_0^{ws} \oplus E_0^{wu}$ of L_0 is smooth, by plaque-expansiveness, [HPS77] guarantees that L_φ as well as its perturbations are dynamically coherent; in particular, $E_F^{ws} \oplus E_F^{wu}$ is uniquely integrable, and by intersecting the leaves of the resulting foliation with the leaves of W_F^s and W_F^u respectively, we obtain foliations integrating E_F^{ws} and E_F^{wu} .

Lemma 2.1. *Let F and G be as above. Then the weak stable distribution E_G^{ws} and weak unstable distribution E_G^{wu} are uniquely integrable.*

Lemma 2.2. *Let F and G be as above and let h be the conjugacy, $h \circ F = G \circ h$. Then $h(W_F^{ws}) = W_G^{ws}$ and $h(W_F^{wu}) = W_G^{wu}$.*

The proofs of these lemmata are quite standard and we only briefly indicate the argument. For more details the reader could consult [GG08], where similar arguments are carried out in detail. In the following, we abbreviate as $W_\varphi^* := W_{L_\varphi}^*$, $*$ $\in \{s, ss, ws, wu, uu, u\}$, the invariant foliations of the de la Llave diffeomorphism L_φ .

By the main result of [DG24] the diffeomorphism G also admits an invariant dominated splitting $E_G^{ss} \oplus E_G^{ws} \oplus E_G^{wu} \oplus E_G^{uu}$. Moreover, the exponential contraction and expansion rates along these subbundles are confined to small intervals $(\lambda - \varepsilon, \lambda + \varepsilon)$, $(\mu - \varepsilon, \mu + \varepsilon)$, $(\mu^{-1} - \varepsilon, \mu^{-1} + \varepsilon)$ and $(\lambda^{-1} - \varepsilon, \lambda^{-1} + \varepsilon)$, respectively.

Denote by h_G the conjugacy to L_φ , $h_G \circ L_\varphi = G \circ h_G$. If E_G^{wu} does not integrate uniquely it still has one-dimensional integral curves γ^{wu} with $\dot{\gamma}^{wu} \in E_G^{wu}$ inside the leaves W_G^u because E_G^{wu} is a one-dimensional distribution. For any such curve γ^{wu} through a point x we have $h_G(\gamma^{wu}) \subset W_\varphi^{wu}(h_G(x))$. This fact follows easily from the following observations:

- (1) $\text{length}(G^n(\gamma^{wu})) \leq (\mu^{-1} + \varepsilon)^n \text{length}(\gamma^{wu})$;
- (2) the lift of the conjugacy h_G to $\mathbb{R}^4$ is a bounded distance away from $\text{id}_{\mathbb{R}^4}$;
- (3) if $p \in W_\varphi^u(q)$ and $p \notin W_\varphi^{wu}(q)$ then $d_{W_\varphi^u}(L_\varphi^n(p), L_\varphi^n(q)) \asymp \lambda^{-n}$, $n \rightarrow \infty$.

If E_G^{wu} is not uniquely integrable then it admits branching integral curves γ^{wu} through a point x all of which must collapse to the same leaf of W_φ^{wu} : $h_G(\gamma^{wu}) \subset W_\varphi^{wu}(x)$, contradicting the fact that h_G is a homeomorphism. This proves Lemma 2.1 (an analogous argument by iterating backward proves the unique integrability of E_G^{ws}).

To establish Lemma 2.2 notice that we proved that $h_G(W_G^{wu}) = W_\varphi^{wu}$. The same argument also gives $h_F(W_F^{wu}) = W_\varphi^{wu}$, where h_F is the conjugacy between F and L_φ . Hence, $h(W_F^{wu}) = h_G^{-1} \circ h_F(W_F^{wu}) = h_G^{-1}(W_\varphi^{wu}) = W_G^{wu}$.

2.1. Geometric mechanism of de la Llave examples. Observe that if two Anosov diffeomorphisms are C^1 conjugate then the differential of the conjugacy maps the fine dominated splitting of the first diffeomorphism to a fine dominated splitting of the second diffeomorphism. Consequently, the conjugacy maps the “web” of invariant foliations of the first diffeomorphisms to that of the second.

Accordingly, if one would like to establish C^1 regularity of the conjugacy it is natural to establish matching of the “webs” of foliations first. In fact, in the context of de la Llave examples, matching of the strong foliations is the main step to obtain rigidity.

Lemma 2.3. *Let L_φ be a de la Llave diffeomorphism. Then there exists a C^1 -neighborhood $\mathcal{U} \ni L_\varphi$ such that if $F \in \mathcal{U}$ and G is conjugate to F via h with the same periodic data and such that $h(W_F^{ss}) = W_G^{ss}$, $h(W_F^{uu}) = W_G^{uu}$, then h is C^{1+H} regular.*

For a detailed proof we refer to [Gog08, Theorem D]. The setting in [Gog08] is more restrictive (only a neighborhood of the linear automorphism is considered), however, given Lemma 2.2 the same proof as in [Gog08] gives Lemma 2.3. Let us simply recall the main steps in the proof of such result:

- for $*$ $\in \{ss, ws, wu, uu\}$, there exist affine structures on the foliation W_F^* , and accordingly for the diffeomorphism G ; more precisely, there exists a continuous family $\{\Psi_{F,x}^*\}_{x \in M}$ of one-dimensional normal forms which linearize the dynamics along the leaves of W_F^* , i.e., for each $x \in M$, $\Psi_{F,x}^*: W_F^*(x) \rightarrow \mathbb{R}$ is a C^{1+H} diffeomorphism⁵ such that

$$\Psi_{F,x}^* \circ F = \|DF(x)|_{E_F^*}\| \cdot \Psi_{F,F(x)};$$

⁵In fact, $\Psi_{F,x}^*$ is smooth for $*$ $\in \{ss, uu\}$; the lack of smoothness of $\Psi_{F,x}^*$ for $*$ $\in \{ws, wu\}$ is due to the lack of regularity of the leaves of W_F^* , which are only C^{1+H} regular in general.

in fact, for $* \in \{wu, uu\}$, $\Psi_{F,x}^*$ is defined as (the expression for $* \in \{ss, ws\}$ is similar, reversing future and past)

$$\Psi_{F,x}^*(y) := \int_x^y \rho_F^*(x, z) dm_{W_F^*(x)}(z), \quad \forall y \in W_F^*(x),$$

where the above integral is taken over the piece of $W_F^*(x)$ from x to y , $dm_{W_F^*(x)}$ being the induced volume on the leaf $W_F^*(x)$, and where

$$\rho_F^*(x, z) := \prod_{k=1}^{+\infty} \frac{\|DF(F^{-k}(x))|_{E_F^*}\|}{\|DF(F^{-k}(z))|_{E_F^*}\|}, \quad \forall z \in W_F^*(x);$$

- if F, G are isospectral, then by Livshits theorem, the functions $\|DF|_{E_F^*}\|$ and $\|DG \circ h|_{E_F^*}\|$ are cohomologous, i.e., for some Hölder continuous function $\gamma^*: M \rightarrow \mathbb{R}$, we have

$$\|DG(h(x))|_{E_F^*}\| = e^{\gamma^*(F(x)) - \gamma^*(x)} \|DF(x)|_{E_F^*}\|, \quad \forall x \in M;$$

if, moreover, the conjugacy h sends W_F^* to W_G^* (it is always true for $* \in \{ws, wu\}$), then for any $x \in M$, $z \in W_F^*(x)$, we have

$$\rho_G^*(h(x), h(z)) = e^{\gamma^*(x) - \gamma^*(z)} \rho_F^*(x, z);$$

then, using the affine structures on W_F^* , it is possible to show that the restriction of h to the leaves of W_F^* is uniformly Lipschitz continuous;

- following the construction of Pesin-Sinai [PS82], [GG08]-[Gog08] construct an F -invariant measure μ whose conditionals along the leaves of W_F^* are absolutely continuous with respect to the induced volume, where for μ -a.e. $x \in M$, the density along $W_F^*(x)$ is proportional to $\rho_F^*(x, \cdot)$; although it is not clear that such measure μ is ergodic, it is shown in [GG08]-[Gog08] that the set of transitive points has full μ -measure, from which it is possible to upgrade the regularity of $h|_{W_F^*(x)}$ from Lipschitz to C^{1+H} ;
- for $* \in \{s, u\}$, the conjugacy h is C^{1+H} along the leaves of the two transverse sub-foliations of W_F^* , hence by Journé's regularity lemma, h is C^{1+H} along the leaves of W_F^* ; since the two foliations W_F^s and W_F^u are transverse, again by Journé's lemma, we conclude that h is C^{1+H} .

Let $p = F^n(p)$ be a periodic point of $F \in \mathcal{U}$ and let $\hat{\mu}_p < \mu_p < \lambda_p < \hat{\lambda}_p$. We say that p is *center-expanding* if $\mu_p \lambda_p > 1$ and *center-contracting* if $\mu_p \lambda_p < 1$. We proceed to state our main technical result which then would supply the hypothesis for Lemma 2.3, which would, in turn, imply the main theorem.

Given a periodic point p_0 we denote by $p_s = F_s^n(p_s) = h_s(p_0)$ its continuation along the family. Also p_0 admits a unique continuation to the neighborhood $\mathcal{U}$ and we will write p_0^G for the corresponding periodic point of $G \in \mathcal{U}$.

Proposition 2.4. *If $\bar{F} \in \mathcal{U}$ and p_0 is a center-expanding periodic point then there exist an arbitrarily C^∞ -small perturbation $\hat{F}$ of $\bar{F}$ and a C^2 -open neighborhood $\mathcal{W}$, $\hat{F} \in \mathcal{W} \subset \mathcal{U}$ such that for all $G \in \mathcal{W}$ and all isospectral deformations $\{G_s\}_{s \in [0,1]}$ of G we have $h_s(W_{G_s, \text{loc}}^{ss}(p_0^G)) = W_{G_s, \text{loc}}^{ss}(p_s^G)$ for all $s \in [0, 1]$.*

Remark 2.5. By reversing the time we also have the symmetric statement for center-contracting periodic point q_0 with the conclusion $h_s(W_{G_s, \text{loc}}^{uu}(p_0^G)) = W_{G_s, \text{loc}}^{uu}(p_s^G)$.

We have the following perturbative counterpart of Proposition 2.4:

Proposition 2.6. *If $\bar{F} \in \mathcal{U}$ and p_0 is a center-expanding periodic point then there exist an arbitrarily C^∞ -small perturbation $\hat{F}$ of $\bar{F}$ and a C^2 -open neighborhood $\mathcal{W}$, $\hat{F} \in \mathcal{W} \subset \mathcal{U}$ such that for all $G \in \mathcal{W}$ and any isospectral diffeomorphism G_1 which is sufficiently C^2 -close to G , the conjugacy h between them satisfies $h(W_{G, \text{loc}}^{ss}(p^G)) = W_{G_1, \text{loc}}^{ss}(p^{G_1})$.*

Proof of Theorem A assuming Proposition 2.4. The following lemma is the starting point.

Lemma 2.7. *Any $F \in \mathcal{U}$ admits arbitrarily C^∞ -small perturbation $\bar{F}$ such that $\bar{F}$ has a center-expanding periodic point p_0 and a center-contracting periodic point, whose eigenvalues satisfy non-resonant conditions (2.1).*

The proof is very standard via a localized perturbation in the neighborhood of periodic points and we omit it.

We can apply Proposition 2.4 to $\bar{F}$ and p_0 given by the lemma to obtain a perturbation $\hat{F}$ and a C^2 -open set $\mathcal{W}(p_0) \subset \mathcal{U}$, $\hat{F} \in \mathcal{W}(p_0)$, such that for all diffeomorphisms $G \in \mathcal{W}(p_0)$ and all isospectral deformations G_s , $s \in [0, 1]$, we have $h_s(W_{G_s, \text{loc}}^{ss}(p_0^G)) = W_{G_s, \text{loc}}^{ss}(p_s^G)$. Further, by taking a smaller perturbation if necessary, we can also assume that $\hat{F}$ still has a center-contracting periodic point which we denote by q_0 .⁶ We can now apply Proposition 2.4 again to $\hat{F}$ and q_0 to obtain a perturbation $\tilde{F}$ and a C^2 -open set $\mathcal{W}(q_0)$, $\tilde{F} \in \mathcal{W}(q_0)$ such that for all $G \in \mathcal{W}(q_0)$ all isospectral deformations G_s , $s \in [0, 1]$ we have $h_s(W_{G_s, \text{loc}}^{uu}(q_0^G)) = W_{G_s, \text{loc}}^{uu}(q_s^G)$, where q_s^G is the continuation of q_0^G . Also, note that by choosing the second perturbation to be sufficiently small we can make sure that $\tilde{F} \in \mathcal{W}(p_0)$ so that $\mathcal{W}(F, \bar{F}, \hat{F}, \tilde{F}) := \mathcal{W}(p_0) \cap \mathcal{W}(q_0)$ is a non-empty set. Note that for any $G \in \mathcal{W}(F, \bar{F}, \hat{F}, \tilde{F})$ and any isospectral deformation G_s we have both $h_s(W_{G_s, \text{loc}}^{uu}(q_0^G)) = W_{G_s, \text{loc}}^{uu}(q_s^G)$ and $h_s(W_{G_s, \text{loc}}^{ss}(p_0^G)) = W_{G_s, \text{loc}}^{ss}(p_s^G)$.

The open set $\mathcal{W}(F, \bar{F}, \hat{F}, \tilde{F})$ is near F , but F does not belong to its closure. To fix this problem we do the following. For every $k \geq 1$ we can repeat the same construction while ensuring that $d_{C^\infty}(F, \bar{F}_k) < 1/k$, $d_{C^\infty}(\bar{F}_k, \hat{F}_k) < 1/k$ and $d_{C^\infty}(\hat{F}_k, \tilde{F}_k) < 1/k$. Then

$$\mathcal{W}(F) = \bigcup_{k \geq 1} \mathcal{W}(F, \bar{F}_k, \hat{F}_k, \tilde{F}_k)$$

is a C^2 -open set which contains F in its C^∞ -closure. Finally, we let

$$\mathcal{V} = \bigcup_{F \in \mathcal{U}} \mathcal{W}(F).$$

This is the C^2 -open and C^∞ -dense set posited in Theorem A. Indeed let $G \in \mathcal{V}$ and let G_s be an isospectral deformation. Since $G \in \mathcal{W}(F, \bar{F}_k, \hat{F}_k, \tilde{F}_k)$ for some F we have some periodic points p_0^G and q_0^G such that the local strong stable manifolds at p_s^G and local strong unstable manifolds at q_s^G match under the conjugacies. In fact, since we can iterate both backward and forward and the base-points are periodic we immediately deduce matching of global strong manifolds — $h_s(W_G^{ss}(p_0^G)) = W_{G_s}^{ss}(p_s^G)$ and $h_s(W_G^{uu}(q_0^G)) = W_{G_s}^{uu}(q_s^G)$. Now by [GG08, Section 4.4] we also have that $h_s(W_G^{ss}(x)) = W_{G_s}^{ss}(h_s(x))$ for all $x \in W^s(p_0^G)$. Since the full stable manifold $W^s(p_0^G)$ is dense in $\mathbb{T}^4$ we deduce that $h_s(W_G^{ss}) = W_{G_s}^{ss}$ and, similarly, $h_s(W_G^{uu}) = W_{G_s}^{uu}$.

Finally, we can apply Lemma 2.3 and conclude that h_s is C^{1+H} regular. $\square$

We now have established Theorem A modulo the proof of Proposition 2.4. By a very similar argument, we can show Theorem B modulo the proof of Proposition 2.6. We proceed with some more preparations and then the proof of Propositions 2.4-2.6 in the next section.

2.2. Passing to suspension flows. Let $F \in \mathcal{U}$ and let F_s , $s \in [0, 1]$, be an isospectral deformation of $F_0 = F$. Denote by $J_s: \mathbb{T}^4 \rightarrow \mathbb{R}$ the (full) Jacobian of F_s . We pick a sufficiently large constant K such that $K + \log J_s > 0$ for all $s \in [0, 1]$, and consider the suspension flows $X_s^t: M \rightarrow M$ of F_s with the roof function $K + \log J_s$. The suspensions flow inherit the dominated splitting structure. That is, for all $s \in [0, 1]$, and $X^t = X_s^t$ we have an DX^t -invariant splitting $TM = E_X^{ss} \oplus E_X^{ws} \oplus \mathbb{R}X \oplus E_X^{wu} \oplus E_X^{uu}$, and we denote the corresponding foliations by W_X^* . We also denote by $E_X^s := E_X^{ss} \oplus E_X^{ws}$ (resp. $E_X^u := E_X^{wu} \oplus E_X^{uu}$) the full

⁶In fact, the construction of perturbation $\hat{F}$ is localized and the eigenvalues of the center-expanding periodic point won't change under the perturbation.

2-dimensional stable (resp. unstable) distribution of the flow X^t . In the following, we will sometimes drop the index X and write E^* , W^* instead of E_X^* , W_X^* , $*$ $\in \{ss, ws, wu, uu\}$.

Since the deformation is isospectral, we have that the sum over the period of $\log J_s$ is independent of s for any periodic point of F_s . Hence, by the Livshits theorem, we have that flows X_s^t are conjugate to the initial flow X_0^t . We denote the conjugacy by H_s — $H_s \circ X_0^t = X_s^t \circ H_s$, $t \in \mathbb{R}$. (The conjugacies H_s are unique up to composition with the flow and we can pick a continuous family of conjugacies H_s , $s \in [0, 1]$.)

2.3. Local coordinates and templates. We consider the deformation X_s^t , $s \in [0, 1]$, and a periodic point $p_0 = X_0^T(p_0)$ with its continuation p_s . Note that the period T is shared by all points p_s . Denote by $\hat{\mu} < \mu < 1 < \lambda < \hat{\lambda}$ the eigenvalues of $D_{p_s}X_s^T$. Note that all these numbers are independent of $s \in [0, 1]$ since the deformation is assumed to be isospectral. From now on, we will focus on various local considerations in the neighborhood of p_s . All constructions that will follow as well as the various choices that will be made can be made continuous in s , since the hyperbolic structures depend on the dynamics continuously. In order to make the notation lighter, we will omit the index s and simply write X^t for X_s^t , p for p_s , etc. Only in the last step of the proof will we reintroduce dependence on s into the notation as the presence of a family becomes important.

Let Σ_p be a smooth 4-dimensional local transversal to the flow which contains the local stable manifold $W_{\text{loc}}^s(p)$ and the local unstable manifold $W_{\text{loc}}^u(p)$. Let us denote by $\Lambda = (\lambda_i)_{i=1}^4 := (\hat{\mu}, \mu, \lambda, \hat{\lambda}) \in \mathbb{R}^4$ the vector of eigenvalues; we can assume that there are no resonances of the form

$$(2.1) \quad \lambda_i = \Lambda^\alpha, \quad \forall i \in \{1, \dots, 4\}, \quad \forall \text{ multi-indices } \alpha \in \mathbb{Z}^4, |\alpha| \leq 3.$$

(Indeed, the resonances at a periodic point can be perturbed away and the absence of resonances is an open property.) Then, by Sternberg linearization theorem, there exists a C^3 parametrization $\Phi_p: (-1, 1)^4 \rightarrow \Sigma_p$ such that the Poincaré return map Π_p to Σ_p is linearized,

$$(2.2) \quad \mathcal{L} := \Phi_p^{-1} \circ \Pi_p \circ \Phi_p: (\hat{\xi}, \xi, \eta, \hat{\eta}) \mapsto (\hat{\mu}\hat{\xi}, \mu\xi, \lambda\eta, \hat{\lambda}\hat{\eta}).$$

The parametrization Φ_p can be easily extended to an actual chart $\iota_p: (-1, 1)^5 \rightarrow M$ around $p = \iota(0, 0, 0, 0, 0)$ such that $\Sigma_p = \iota_p((-1, 1)^2 \times \{0\} \times (-1, 1)^2)$, and for any $(\hat{\xi}, \xi, t, \eta, \hat{\eta}) \in (-1, 1)^5$,

- (1) $\iota_p(\hat{\xi}, \xi, 0, \eta, \hat{\eta}) = \Phi_p(\hat{\xi}, \xi, \eta, \hat{\eta})$;
- (2) $\iota_p(\hat{\xi}, \xi, t, \eta, \hat{\eta}) = X^t(\iota_p(\hat{\xi}, \xi, 0, \eta, \hat{\eta}))$.

Note that the local weak and strong stable/unstable manifolds through p are just the coordinate axes in this parametrization.

Using this chart we define two functions, the *weak stable template* $\mathcal{T}_p^{ws}$ and the *strong stable template* $\mathcal{T}_p^{ss}$, which represent the angular coordinates of (the image in normal coordinates of) the 2-plane $E_X^s = E_X^{ss} \oplus E_X^{ws}$ as we move along the local unstable manifold $W_{\text{loc}}^u(p)$:

$$(2.3) \quad \begin{aligned} D\iota_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(0, 1, \mathcal{T}_p^{ws}(\eta, \hat{\eta}), *, *)) &\subset E_X^s(\Phi_p(0, 0, \eta, \hat{\eta})), \\ D\iota_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(1, 0, \mathcal{T}_p^{ss}(\eta, \hat{\eta}), *, *)) &\subset E_X^s(\Phi_p(0, 0, \eta, \hat{\eta})). \end{aligned}$$

Remark 2.8. The subspace $D\iota_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(0, 1, \mathcal{T}_p^{ws}(\eta, \hat{\eta}), *, *))$ is close to $E_X^{ws}(\Phi_p(0, 0, \eta, \hat{\eta}))$ and $D\iota_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(1, 0, \mathcal{T}_p^{ss}(\eta, \hat{\eta}), *, *))$ is close to $E_X^{ss}(\Phi_p(0, 0, \eta, \hat{\eta}))$, hence the choice of notation. In fact, $\mathcal{T}_p^{ws}$, resp. $\mathcal{T}_p^{wu}$ can be thought of as the temporal coordinate of the bundle $E_X^{ws} \oplus E_X^u$, resp. $E_X^s \oplus E_X^{wu}$, along $W_{\text{loc}}^u(p)$, resp. $W_{\text{loc}}^s(p)$, in connection with items (1)-(2) in Theorem C.

We also denote by $\tau_p: \Sigma_p \rightarrow \mathbb{R}$ the first return time of the flow X^t on the transversal Σ_p , specifically, for $x \in \Sigma_p$, $\Pi_p(x) = X^{\tau_p(x)}(x)$. In the following, we identify τ_p with $\tau_p \circ \Phi_p$. Note that by construction

$$(2.4) \quad \tau_p(\cdot, \cdot, 0, 0) = \tau_p(0, 0, \cdot, \cdot) \equiv T.$$

2.4. Coarse local coordinates. While the local sections and charts we have described above are very nice for performing calculations, they do have one drawback: since the linearization procedure depends on higher jets at the fixed point, the sections vary continuously with respect to the flow only in a sufficiently high topology. Meanwhile we have a family of suspension flows which is continuous in C^1 topology. Since continuity in parameter s will play an important role, we will also need to have another family of transversals and charts which would vary continuously in C^1 topology with s .

Specifically, we still consider transversals Σ_p , but impose less rigid conditions on the parametrization $\hat{\Phi}_p: (-1, 1)^4 \rightarrow \Sigma_p$. Namely, we only require the following weaker conditions.

- (1) $\hat{\Phi}_p((-1, 1) \times (-1, 1) \times \{0\} \times \{0\}) = W_{X, \text{loc}}^s(p);$
- (2) $\hat{\Phi}_p((-1, 1) \times \{0\} \times \{0\} \times \{0\}) = W_{X, \text{loc}}^{ss}(p);$
- (3) $D\hat{\Phi}_p(\frac{\partial}{\partial \xi})(p) = E_X^{ws}(p);$
- (4) $\hat{\Phi}_p(\{0\} \times \{0\} \times (-1, 1) \times (-1, 1)) = W_{X, \text{loc}}^u(p);$
- (5) $\hat{\Phi}_p(\{0\} \times \{0\} \times \{0\} \times (-1, 1)) = W_{X, \text{loc}}^{uu}(p);$
- (6) $D\hat{\Phi}_p(\frac{\partial}{\partial \eta})(p) = E_X^{wu}(p).$

We denote by $\hat{\Pi}_p$ the corresponding Poincaré return map, $\hat{\Pi}_p := \hat{\Phi}_p^{-1} \circ \Pi_p \circ \hat{\Phi}_p$. By the above conditions we have $D\hat{\Pi}_p(0, 0, 0, 0) = \mathcal{L}$. We extend the parametrization $\hat{\Phi}_p$ to an actual local chart $\hat{i}_p$ in exactly the same way as before and we will need the corresponding return time $\hat{\tau}_p = \tau \circ \hat{\Phi}_p$. Finally we also define corresponding templates in the exact same way:

$$\begin{aligned} D\hat{i}_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(0, 1, \hat{\mathcal{T}}_p^{ws}(\eta, \hat{\eta}), *, *)) &\subset E_X^s(\hat{\Phi}_p(0, 0, \eta, \hat{\eta})), \\ D\hat{i}_p(0, 0, 0, \eta, \hat{\eta})(\text{Span}(1, 0, \hat{\mathcal{T}}_p^{ss}(\eta, \hat{\eta}), *, *)) &\subset E_X^s(\hat{\Phi}_p(0, 0, \eta, \hat{\eta})). \end{aligned}$$

Remark 2.9. It will be important that for the C^1 family of suspension flows X_s^t (which comes from the C^2 family of Anosov diffeomorphisms F_s) the coarse charts at p_s , $s \in [0, 1]$, can be chosen in a continuous manner. Indeed, this is clear because all local invariant manifolds at through p_s vary continuously in C^1 topology on the flows and all that is required of the coarse charts $\hat{\Phi}_{p_s}$ is that they align in certain way with these invariant manifolds.

In fact the proof simplifies quite a bit for smooth families, since then for the most part the reader can pretend that coarse charts are the same as linearizing charts. Indeed, the main point of coarse charts is that they depend continuously on the parameter and, if the family is smooth, the linearizing charts do depend continuously on the parameter. Still coarse charts are needed for Lemma 3.12.

2.5. Prescribed families of shadowing periodic orbits. As above, we consider a center-expanding periodic point $p \in M$, of period $T > 0$, with eigenvalues $\hat{\mu} < \mu < 1 < \lambda < \hat{\lambda}$.

We fix some homoclinic point $q \in W_{\text{loc}}^u(p)$. We fix a time $T' > 0$ with the property that $q' = X^{T'}(q) \in W_{\text{loc}}^s(p)$. Without loss of generality, we can assume that $q, q' \in \Sigma_p$. The following shadowing result is standard; see e.g. [GLH25, Lemma 4.1] for more details.

Lemma 2.10. *There exist a constant $C_0 > 0$ and an integer $n_0 \in \mathbb{N}$ such that for $n \geq n_0$, there exists a unique periodic point $p_n \in \Sigma_p$ of period*

$$T_n \simeq nT + T'$$

such that

$$d(X^t(p_n), X^t(q)) \leq C_0 \mu^{\frac{n}{2}}, \quad \forall t \in \left[\frac{-nT}{2}, \frac{nT}{2} + T' \right].$$

3. THE PROOF

This section is devoted to the proof of Proposition 2.4. Recall that we have reduced the proof of the main theorem to Proposition 2.4.

3.1. Asymptotic formula for full Jacobian along shadowing periodic orbits (uniform in the family). Let $p = X^T(p)$ be a center-expanding periodic point. Let $q = \Phi_p(0, 0, \eta_\infty, \hat{\eta}_\infty) \in W_{\text{loc}}^u(p)$ be a homoclinic point, and fix $T' > 0$ such that $q' = X^{T'}(q) = \Phi_p(\hat{\xi}_\infty, \xi_\infty, 0, 0) \in W_{\text{loc}}^s(p)$. Let $(p_n = \Phi_p(\hat{\xi}_n, \xi_n, \eta_n, \hat{\eta}_n))_n$ be the sequence of shadowing periodic points associated to q given by Lemma 2.10, and let $p'_n = X^{\tau_n}(p_n) = \Phi_p(\hat{\xi}'_n, \xi'_n, \eta'_n, \hat{\eta}'_n) \in \Sigma_p$ the point in the orbit of p_n which is closest to q' . In particular, $\tau_n \approx T'$, $p_n \approx X^{nT}(p'_n)$, and by the shadowing lemma,

$$(3.1) \quad (\hat{\xi}_n, \xi_n, \eta_n, \hat{\eta}_n) = (0, 0, \eta_\infty, \hat{\eta}_\infty) + O(\mu^{\frac{n}{2}}), \quad (\hat{\xi}'_n, \xi'_n, \eta'_n, \hat{\eta}'_n) = (\hat{\xi}_\infty, \xi_\infty, 0, 0) + O(\mu^{\frac{n}{2}}).$$

Recall that we denote by $T_n \approx nT + T'$ the period of the point p_n .

Proposition 3.1. *Given a center-expanding periodic point $p = X^T(p)$, any homoclinic point $q \in W_{\text{loc}}^u(p)$, $q' = X^{T'}(q) \in W_{\text{loc}}^s(p)$, and the corresponding sequence of shadowing periodic points (p_n) associated to q, q' , their periods T_n obey the following asymptotic expansion:*

$$T_n = nT + T' + \xi_\infty (\mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_p(\eta_\infty, \hat{\eta}_\infty)) \mu^n + O(\theta^n),$$

with $\theta := \mu^{\frac{2 \log \lambda}{\log \lambda - \log \mu}} \in (\mu^2, \mu)$, and

$$P_p(\eta_\infty, \hat{\eta}_\infty) = - \sum_{\ell=1}^{+\infty} \mu^{-\ell} \partial_2 \tau_p(0, 0, \lambda^{-\ell} \eta_\infty, \hat{\lambda}^{-\ell} \hat{\eta}_\infty),$$

where ∂_2 is the partial derivative with respect to the second variable ξ .

In the following, we abuse notation and denote $\mathcal{T}_p^{ws}(q) := \mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty)$, $P_p(q) := P_p(\eta_\infty, \hat{\eta}_\infty)$.

Remark 3.2. Since the conjugacy map Φ_p between Π_p and its linearization is C^3 , after this change of coordinates the first return time τ_p is also C^3 . Moreover, $\partial_2 \tau_p(0, 0, 0, 0) = 0$, hence by Taylor expansion, we see that the terms $P_p^{(\ell)}(\eta_\infty, \hat{\eta}_\infty) := \mu^{-\ell} \partial_2 \tau_p(0, 0, \lambda^{-\ell} \eta_\infty, \hat{\lambda}^{-\ell} \hat{\eta}_\infty)$ of the series defining $P_p(\eta_\infty, \hat{\eta}_\infty)$ decay at least like $(\mu\lambda)^{-\ell}$ (recall $(\mu\lambda)^{-1} < 1$ since p is center-expanding), so $(\eta_\infty, \hat{\eta}_\infty) \mapsto P_p(\eta_\infty, \hat{\eta}_\infty)$ is indeed a well-defined continuous function. Moreover, for any $\ell \in \mathbb{N}$, the function $(\eta_\infty, \hat{\eta}_\infty) \mapsto P_p^{(\ell)}(\eta_\infty, \hat{\eta}_\infty)$ is C^1 (in fact, C^2), and both partial derivatives of $P_p^{(\ell)}$ decay faster than $(\mu\lambda)^{-\ell}$, which ensures that the function P_p is, in fact, C^1 regular.

Lemma 3.3. *We have the following expansions*

$$\begin{aligned} (\hat{\xi}_n, \xi_n, \eta_n, \hat{\eta}_n) - (0, 0, \eta_\infty, \hat{\eta}_\infty) &= (\hat{\mu}^n \hat{\xi}_\infty + O(\hat{\mu}^n \mu^n), \mu^n \xi_\infty + O(\mu^{2n}), O(\lambda^{-n}), O(\lambda^{-n})), \\ (\hat{\xi}'_n, \xi'_n, \eta'_n, \hat{\eta}'_n) - (\hat{\xi}_\infty, \xi_\infty, 0, 0) &= (O(\mu^n), O(\mu^n), \lambda^{-n} \eta_\infty + O(\lambda^{-2n}), \hat{\lambda}^{-n} \hat{\eta}_\infty + O(\hat{\lambda}^{-n} \lambda^{-n})). \end{aligned}$$

Proof. By (3.1), and since in the chart Φ_p the dynamics is given by the linear map $\mathcal{L}$ (2.2), we have

$$(3.2) \quad (\eta'_n, \hat{\eta}'_n) = (\lambda^{-n} \eta_n, \hat{\lambda}^{-n} \hat{\eta}_n) = (\lambda^{-n} \eta_\infty + O(\lambda^{-n} \mu^{\frac{n}{2}}), \hat{\lambda}^{-n} \hat{\eta}_\infty + O(\hat{\lambda}^{-n} \mu^{\frac{n}{2}})),$$

$$(3.3) \quad (\hat{\xi}_n, \xi_n) = (\hat{\mu}^n \hat{\xi}'_n, \mu^n \xi'_n) = (\hat{\mu}^n \hat{\xi}_\infty + O(\hat{\mu}^n \mu^{\frac{1}{2}n}), \mu^n \xi_\infty + O(\mu^{\frac{3}{2}n})).$$

Let us denote by $\bar{\Pi} = X^{\bar{\tau}}: U_q \rightarrow U_{q'}$ the Poincaré map from a neighborhood $U_q \subset \Sigma_p$ of q to a neighborhood $U_{q'} \subset \Sigma_p$ of q' , with $\bar{\tau} \approx T'$. Since $\Phi_p^{-1}(W_{\text{loc}}^s(q')) \subset (-1, 1)^2 \times \{(0, 0)\}$ and $\Phi_p^{-1}(W_{\text{loc}}^u(q'))$ are transverse, estimate (3.2) shows that the unstable distance $d_{W^u}(q', p'_n)$ between q' and p'_n is of order $O(\lambda^{-n})$. Applying the C^3 diffeomorphism $\bar{\Pi}^{-1}$, we see that the unstable distance $d_{W^u}(q, p_n)$ between the points $q = \bar{\Pi}^{-1}(q')$ and $p_n = \bar{\Pi}^{-1}(p'_n)$ is also of order $O(\lambda^{-n})$. Since $\Phi_p^{-1}(W_{\text{loc}}^u(q)) \subset \{(0, 0)\} \times (-1, 1)^2$, we deduce that

$$(\eta_n, \hat{\eta}_n) = (\eta_\infty, \hat{\eta}_\infty) + O(\lambda^{-n}).$$

Using estimate (3.3), analogous reasoning leads to

$$(\hat{\xi}'_n, \xi'_n) = (\hat{\xi}_\infty, \xi_\infty) + O(\mu^n).$$

Applying $\mathcal{L}^{\pm n}$ to the last two estimates as in (3.2)-(3.3) allows us to obtain upgraded estimates

$$\begin{aligned}(\hat{\xi}_n, \xi_n) &= (\hat{\mu}^n \hat{\xi}_\infty + O(\hat{\mu}^n \mu^n), \mu^n \xi_\infty + O(\mu^{2n})), \\(\eta'_n, \hat{\eta}'_n) &= (\lambda^{-n} \eta_\infty + O(\lambda^{-2n}), \hat{\lambda}^{-n} \hat{\eta}_\infty + O(\hat{\lambda}^{-n} \lambda^{-n})),\end{aligned}$$

which concludes the proof. $\square$

Again, since the dynamics in the charts is linear (2.2), the previous estimates at the entrance/exit points p'_n, p_n can be immediately propagated to the whole orbit:

Corollary 3.4. *For any $\ell \in \{0, \dots, n\}$, let*

$$\begin{aligned}\tilde{p}_n(-\ell) &:= \Phi_p^{-1}(\Pi_p^{-\ell}(p_n)) = (\hat{\mu}^{-\ell} \hat{\xi}_n, \mu^{-\ell} \xi_n, \lambda^{-\ell} \eta_n, \hat{\lambda}^{-\ell} \hat{\eta}_n), \\ \tilde{q}(-\ell) &:= \Phi_p^{-1}(\Pi_p^{-\ell}(q)) = (0, 0, \lambda^{-\ell} \eta_\infty, \hat{\lambda}^{-\ell} \hat{\eta}_\infty), \\ \tilde{q}'(\ell) &:= \Phi_p^{-1}(\Pi_p^\ell(q')) = (\hat{\mu}^\ell \hat{\xi}_\infty, \mu^\ell \xi_\infty, 0, 0).\end{aligned}$$

Then, we have

$$\begin{aligned}\tilde{p}_n(-\ell) - \tilde{q}(-\ell) &= (\hat{\mu}^{-\ell} \hat{\xi}_n, \mu^{-\ell} \xi_n, \lambda^{-\ell} \eta_n, \hat{\lambda}^{-\ell} \hat{\eta}_n) - (0, 0, \lambda^{-\ell} \eta_\infty, \hat{\lambda}^{-\ell} \hat{\eta}_\infty) \\ &= (\hat{\mu}^{n-\ell} \hat{\xi}_\infty + O(\mu^n \hat{\mu}^{n-\ell}), \mu^{n-\ell} \xi_\infty + O(\mu^{2n-\ell}), O(\lambda^{-(n+\ell)}), O(\lambda^{-n} \hat{\lambda}^{-\ell})), \\ \tilde{p}_n(-\ell) - \tilde{q}'(n-\ell) &= (\hat{\mu}^{n-\ell} \hat{\xi}'_n, \mu^{n-\ell} \xi'_n, \lambda^{n-\ell} \eta'_n, \hat{\lambda}^{n-\ell} \hat{\eta}'_n) - (\hat{\mu}^{n-\ell} \hat{\xi}_\infty, \mu^{n-\ell} \xi_\infty, 0, 0) \\ &= (O(\mu^n \hat{\mu}^{n-\ell}), O(\mu^{2n-\ell}), \lambda^{-\ell} \eta_\infty + O(\lambda^{-(n+\ell)}), \hat{\lambda}^{-\ell} \hat{\eta}_\infty + O(\lambda^{-n} \hat{\lambda}^{-\ell})).\end{aligned}$$

Corollary 3.5. *As $n \rightarrow +\infty$, we have the asymptotic expansion*

$$(3.4) \quad \sum_{\ell=1}^n \tau_p(\tilde{p}_n(-\ell)) = nT + \xi_\infty \left[\sum_{\ell=1}^{+\infty} \partial_2 \tau_p(\tilde{q}(-\ell)) \mu^{-\ell} \right] \mu^n + O(\theta^n),$$

where $\theta := \mu^{\frac{2 \log \lambda}{\log \lambda - \log \mu}} \in (\mu^2, \mu)$.

Proof. For any $n \geq 0$, we define the integer

$$(3.5) \quad \ell_n := \left\lfloor \frac{-\log \mu}{\log \lambda - \log \mu} n \right\rfloor, \quad \text{so that } \mu^{n-\ell_n} \approx \lambda^{-\ell_n}.$$

Recall that the roof function τ_p in normal coordinates is C^3 . By (2.4), we have $\partial_3 \tau_p(0, 0, \cdot, \cdot) = \partial_4 \tau_p(0, 0, \cdot, \cdot) \equiv 0$; similarly, $\partial_1 \tau_p(\cdot, \cdot, 0, 0) = \partial_2 \tau_p(\cdot, \cdot, 0, 0) \equiv 0$. Then, by Corollary 3.4, and by Taylor expansion, for $n \gg 1$ and for any $\ell \in \{1, \dots, \ell_n\}$, we have

$$\begin{aligned}\tau_p(\tilde{p}_n(-\ell)) &= \tau_p(\tilde{q}(-\ell)) + \partial_1 \tau_p(\tilde{q}(-\ell))(\hat{\mu}^{n-\ell} \hat{\xi}_\infty + O(\mu^n \hat{\mu}^{n-\ell})) \\ &\quad + \partial_2 \tau_p(\tilde{q}(-\ell))(\mu^{n-\ell} \xi_\infty + O(\mu^{2n-\ell})) + O \left[\sup_{[\tilde{q}(-\ell), \tilde{p}_n(-\ell)]} D^2 \tau_p(\tilde{p}_n(-\ell) - \tilde{q}(-\ell))^2 \right] \\ &= T + \left[\partial_1 \tau_p(0_{\mathbb{R}^4}) + O(\lambda^{-\ell}) \right] (\hat{\mu}^{n-\ell} \hat{\xi}_\infty + O(\mu^n \hat{\mu}^{n-\ell})) \\ &\quad + \partial_2 \tau_p(\tilde{q}(-\ell)) \mu^{n-\ell} \xi_\infty + \left[\partial_2 \tau_p(0_{\mathbb{R}^4}) + O(\lambda^{-\ell}) \right] O(\mu^{2n-\ell}) \\ &\quad + O \left[\left[D^2 \tau_p(0_{\mathbb{R}^4}) + O(\lambda^{-\ell}) \right] (\tilde{p}_n(-\ell) - \tilde{q}(-\ell))^2 \right] \\ &= T + \partial_2 \tau_p(\tilde{q}(-\ell)) \mu^{n-\ell} \xi_\infty + O(\lambda^{-\ell} \hat{\mu}^{n-\ell}) + O(\lambda^{-\ell} \mu^{2n-\ell}) \\ &\quad + O(\mu^{n-\ell} \lambda^{-(n+\ell)}) + O(\lambda^{-\ell} \mu^{2(n-\ell)}).\end{aligned}$$

In the above estimates we have used that:

- for $\ell \in \{1, \dots, \ell_n\}$, the main discrepancy between the periodic point $\tilde{p}_n(-\ell)$ and the homoclinic point $\tilde{q}(-\ell)$ is along the second (ξ) coordinate, and it is of order $\mu^{n-\ell}$;
- $\partial_1 \tau_p(0_{\mathbb{R}^4}) = \partial_2 \tau_p(0_{\mathbb{R}^4}) = 0$;

- the Hessian $D^2\tau_p(0_{\mathbb{R}^4})$ has $\partial_1^2\tau_p(0_{\mathbb{R}^4}) = \partial_1\partial_2\tau_p(0_{\mathbb{R}^4}) = \partial_2^2\tau_p(0_{\mathbb{R}^4}) = 0$, and similarly, $\partial_3^2\tau_p(0_{\mathbb{R}^4}) = \partial_3\partial_4\tau_p(0_{\mathbb{R}^4}) = \partial_4^2\tau_p(0_{\mathbb{R}^4}) = 0$; in other words, its only nonzero terms are those involving a mixed derivative of the form stable direction (ws or ss) vs. unstable direction (wu or uu).

Recall that $\mu\lambda > 1$, $\hat{\mu} < \mu < 1$, and that $\partial_2\tau_p(\tilde{q}(-\ell)) = O(\lambda^{-\ell})$. We deduce that

$$\sum_{\ell=1}^{\ell_n} \tau_p(\tilde{p}_n(-\ell)) = \ell_n T + \xi_\infty \left[\sum_{\ell=1}^{+\infty} \partial_2\tau_p(\tilde{q}(-\ell)) \mu^{-\ell} \right] \mu^n + O((\mu\lambda)^{-\ell_n}) \mu^n.$$

Note that by (3.5), we have

$$(\mu\lambda)^{-\ell_n} \mu^n \approx \lambda^{-2\ell_n} \approx \mu^\gamma, \text{ with } \gamma := \frac{2 \log \lambda}{\log \lambda - \log \mu} \in (1, 2).$$

Moreover, for any $\ell \in \{\ell_n + 1, \dots, n\}$, we have

$$\begin{aligned} \tau_p(\tilde{p}_n(-\ell)) &= \tau_p(\tilde{q}'(n-\ell)) + O \left[\sup_{[\tilde{q}'(n-\ell), \tilde{p}_n(-\ell)]} \partial_3\tau_p \right] \lambda^{-\ell} + O \left[\sup_{[\tilde{q}'(n-\ell), \tilde{p}_n(-\ell)]} \partial_4\tau_p \right] \hat{\lambda}^{-\ell} \\ &= T + O(\mu^{n-\ell} \lambda^{-\ell}). \end{aligned}$$

Gathering together the above estimates, we thus conclude that

$$\sum_{\ell=1}^n \tau_p(\tilde{p}_n(-\ell)) = nT + \xi_\infty \left[\sum_{\ell=1}^{+\infty} \partial_2\tau_p(\tilde{q}(-\ell)) \mu^{-\ell} \right] \mu^n + O(\mu^\gamma). \quad \square$$

Let us now explain how to conclude the proof of Proposition 3.1. In Corollary 3.5, we have computed the total discrepancy between the period T and the return times to the transverse section $\Sigma_p = \{t = 0\}$ of the iterates $\Pi_p^{-\ell}(p_n)$, $\ell \in \{1, \dots, n\}$, which accounts for the term $\xi_\infty P_p(\eta_\infty, \hat{\eta}_\infty) \mu^n$ in the asymptotic expansion of the periods T_n . It thus remains to deal with the time of the excursion from p_n to p'_n , which itself shadows closely the homoclinic excursion from q to q' . In particular, this excursion time is of the form $T' + \bar{\tau}(p_n)$, with $\bar{\tau}(p_n) = o(1)$. In fact, $\bar{\tau}(p_n)$ accounts for the term involving the weak stable template $\mathcal{T}_p^{ws}$ in Proposition 3.1:

Lemma 3.6. *As $n \rightarrow +\infty$, we have the following asymptotic expansion:*

$$(3.6) \quad \bar{\tau}(p_n) = \xi_\infty \mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty) \mu^n + O(\max(\lambda^{-n}, \mu^{2n})).$$

Proof. Let us sketch the proof of (3.6); for more details, we refer the reader to [GLH25, Lemma 4.7-Claim 4.8]). The main steps are as follows:

- we claim that $\bar{\tau}(p_n)$ is equal to the time that the point p_n needs to reach the stable manifold $W_{X, \text{loc}}^s(q)$ (which is *a priori* not contained in Σ_p); indeed, the point $X^{\bar{\tau}(p_n)}(p_n)$ travels together with q , hence reaches Σ_p exactly after time T' ;
- near the homoclinic point q (i.e., for $|\hat{\xi}|, |\xi| \ll 1$ in normal coordinates), the stable manifold $W_{X, \text{loc}}^s(q)$ is well approximated by its tangent space $E_X^s(q)$, which in normal coordinates can be expressed as a graph

$$(3.7) \quad \{(\hat{\xi}, \xi, \hat{\xi} \mathcal{T}_p^{ss}(\eta_\infty, \hat{\eta}_\infty) + \xi \mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty), *, *) : \hat{\xi}, \xi \in \mathbb{R}\},$$

by the definition of the weak-stable and strong stable templates in (2.3);

- by Lemma 3.3, the leading term of the discrepancy between p_n and q is along the weak-stable direction, i.e., the second coordinate in normal coordinates, and it is of order $\xi_\infty \mu^n + O(\lambda^{-n})$;
- since $p_n \in \Sigma_p = \{t = 0\}$, by (3.7), we deduce that

$$\bar{\tau}(p_n) = \xi_\infty \mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty) \mu^n + O(\max(\lambda^{-n}, \mu^{2n})),$$

where the error term μ^{2n} comes from the approximation of $W_{X, \text{loc}}^s(q)$ by its tangent space $E_X^s(q)$. $\square$

Since $T_n = \sum_{\ell=1}^n \tau_p(\tilde{p}_n(-\ell)) + \bar{\tau}(p_n) + T'$, combining the asymptotic expansion of the former sum obtained in Corollary 3.5 with the estimate of $\bar{\tau}(p_n)$ derived in Lemma 3.6, this concludes the proof of Proposition 3.1. $\square$

3.2. Reinterpreting the leading term of the asymptotic formula in the coarse charts. As explained in Subsection 2.4, the computations above are performed in linearizing charts which vary continuously with respect to the flow only in a sufficiently high topology (in fact, C^4 topology). Here, given a center-expanding point $p \in M$, and homoclinic points q, q' as above, we explain how the coefficient by the leading exponentially small term in the asymptotic formula in Proposition 3.1 can be reinterpreted in coarser charts which depend continuously on the flow in C^1 topology. We let

$$\begin{aligned}\zeta_p(q) &:= \mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_p(\eta_\infty, \hat{\eta}_\infty), \\ \omega_p(q, q') &:= \xi_\infty \cdot \zeta_p(q),\end{aligned}$$

so that $\omega_p(q, q')$ is the coefficient by the leading exponentially small term in the asymptotic formula derived in Proposition 3.1. We let $\hat{\Phi}_p$, $\hat{\Pi}_p$, $\hat{\tau}_p$, and $\hat{\mathcal{T}}_p^{ws}$ be the associated objects in the coarse charts introduced in Subsection 2.4. For any $\ell \geq 0$, we let $\hat{\mu}_{p,q}(-\ell)$ be the stable Jacobian of $D\hat{\Pi}_p^{-\ell}$ at $\hat{\Phi}_p^{-1}(q)$, and let

$$\begin{aligned}(0, 0, \eta_\infty^\circ, \hat{\eta}_\infty^\circ) &:= \hat{\Phi}_p^{-1}(q), \\ (\hat{\xi}_\infty^\circ, \xi_\infty^\circ, 0, 0) &:= \hat{\Phi}_p^{-1}(q'), \\ (3.8) \quad \hat{P}_p(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) &:= - \sum_{\ell=1}^{+\infty} \hat{\mu}_{p,q}(-\ell) \partial_2 \hat{\tau}_p(\hat{\Pi}_p^{-\ell}(\hat{\Phi}_p^{-1}(q))),\end{aligned}$$

$$\begin{aligned}\hat{\zeta}_p(q) &:= \hat{\mathcal{T}}_p^{ws}(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) - \hat{P}_p(\eta_\infty^\circ, \hat{\eta}_\infty^\circ), \\ \hat{\omega}_p(q, q') &:= \xi_\infty^\circ \cdot \hat{\zeta}_p(q).\end{aligned}$$

The following result is precisely the content of [GLH25, Lemma 7.16]:

Lemma 3.7. *There exists a positive constant $\vartheta_p(q) > 0$ such that $\zeta_p(q) = \vartheta_p(q) \cdot \hat{\zeta}_p(q)$.*

Remark 3.8. In fact, in the exact same way, we can define two functions $\zeta_p, \hat{\zeta}_p$ along $W_X^u(p)$; moreover, Lemma [GLH25, Lemma 7.16] guarantees that there exists a positive function $\vartheta_p: W_X^u(p) \rightarrow \mathbb{R}_+$ such that $\zeta_p = \vartheta_p \cdot \hat{\zeta}_p$.

Corollary 3.9. *The quantity $\hat{\omega}_p(q, q')$ is > 0 , resp. $= 0$, resp. < 0 if and only if the quantity $\omega_p(q, q')$ is > 0 , resp. $= 0$, resp. < 0 .*

Proof. It follows directly from Lemma 3.7; indeed, ξ_∞ and ξ_∞° have the same sign, since they encode the position of $q' \in W_{X,\text{loc}}^s(p)$ relative to $W_{X,\text{loc}}^{ss}(p)$ in the respective charts $\Phi_p, \hat{\Phi}_p$. $\square$

3.3. Non-vanishing of the leading term after a C^∞ perturbation.

Proposition 3.10. *Fix an Anosov diffeomorphism $\bar{F}$, a center-expanding periodic point p for the suspension flow as above. Fix any $(\eta, \hat{\eta}) \in (-1, 1)^2$ such that $q = \hat{\Phi}_p(0, 0, \eta, \hat{\eta})$ is homoclinic. For any $\varepsilon > 0$, there exists a C^∞ diffeomorphism $\hat{F}$, $d_{C^\infty}(\bar{F}, \hat{F}) < \varepsilon$, and $r > 0$, such that for any C^∞ diffeomorphism G satisfying $d_{C^2}(G, \hat{F}) < r$, denoting by Y^t the suspension flow of G , and by p^G the continuation of p , the functions $\hat{\mathcal{T}}_{p^G}^{ws}$ and $\hat{P}_{p^G}$ associated to Y^t (in the coarse chart $\hat{\Phi}_{p^G}$) satisfy*

$$\hat{\mathcal{T}}_{p^G}^{ws}(\eta', \hat{\eta}') - \hat{P}_{p^G}(\eta', \hat{\eta}') \neq 0, \quad \forall (\eta', \hat{\eta}') \in B((\eta, \hat{\eta}), r).$$

The notation for the diffeomorphisms $\bar{F}$, $\hat{F}$ and G in Proposition 3.10 is consistent with the notation in Proposition 2.4.

Remark 3.11. In particular, for any diffeomorphism G which is sufficiently C^2 -close to $\hat{F}$, for any homoclinic point $q^G = \Phi_p^G(0, 0, \eta_\infty, \hat{\eta}_\infty) = \hat{\Phi}_p^G(0, 0, \eta_\infty^\circ, \hat{\eta}_\infty^\circ)$ with $(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in B((\eta, \hat{\eta}), r)$, Proposition 3.10 ensures that $\hat{\zeta}_p(q) = \hat{\mathcal{T}}_{p^G}^{ws}(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) - \hat{P}_{p^G}(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \neq 0$, hence also, by Lemma 3.7, $\zeta_p(q) = \mathcal{T}_{p^G}^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_{p^G}(\eta_\infty, \hat{\eta}_\infty) \neq 0$. Consequently, in the asymptotic expansion obtained in Proposition 3.1, the coefficient $\omega_p(q, q') = \xi_\infty(\mathcal{T}_{p^G}^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_{p^G}(\eta_\infty, \hat{\eta}_\infty))$ by the leading exponential term is non-zero unless $\xi_\infty = 0$.

We split the proof of this proposition into two lemmata, showing respectively the C^∞ -density and C^2 -openness of the above property.

Lemma 3.12. *There exists a C^∞ -small perturbation $\hat{F}$ of $\bar{F}$ such that the functions $\hat{\mathcal{T}}_{p, \hat{F}}^{ws}$ and $\hat{P}_{p, \hat{F}}$ associated to the suspension flow of $\hat{F}$ (in the coarse chart $\hat{\Phi}_p$) satisfy*

$$\hat{\mathcal{T}}_{p, \hat{F}}^{ws}(\eta, \hat{\eta}) - \hat{P}_{p, \hat{F}}(\eta, \hat{\eta}) \neq 0.$$

Sketch of the proof. We only explain the basic mechanism of non-vanishing without carrying out detailed computations. This is because the perturbative argument is rather standard, in particular it is used to show that strong stable distribution of an Anosov flows is typically not C^1 regular, see, e.g. [GLH25, Lemma 7.12].

The perturbation $\hat{F}$ will be localized at a homoclinic point which corresponds to the preimage of the homoclinic point, i.e., $\hat{\Pi}^{-1}(0, 0, \eta, \hat{\eta})$. Further, the perturbation can be arranged so that the periodic point p and all invariant manifolds through p (and dynamics on them) remain the same for $\hat{F}$. Then we can use exactly the same coarse chart $\hat{\Phi}_p$ for both $\bar{F}$ and $\hat{F}$ because it is only adapted with respect to these invariant manifolds.

The template $\hat{\mathcal{T}}_p^{ws}(\eta, \hat{\eta})$ is defined via the chart and the stable subspace $E_X^s(q) = E_X^s(\hat{\Phi}_p(0, 0, \eta, \hat{\eta}))$. Hence, $\hat{\mathcal{T}}_p^{ws}(\eta, \hat{\eta})$ is completely determined by the future of q and the chart $\hat{\Phi}_p$. Since the forward orbit of q is asymptotic to p , it never intersects a small neighborhood of $\hat{\Phi}_p(\hat{\Pi}^{-1}(0, 0, \eta, \hat{\eta}))$, the stable manifold for such localized perturbation at q remains the same; and since we are still using the same chart, we conclude that the template remains exactly the same for the perturbation $\hat{F}$:

$$\hat{\mathcal{T}}_{p, \hat{F}}^{ws}(\eta, \hat{\eta}) = \hat{\mathcal{T}}_p^{ws}(\eta, \hat{\eta}).$$

Now recall that $\hat{P}_p$ is given by the series (3.8) which go into the past along the local unstable manifold of p . The only term of these series which could change under such localized perturbation is

$$\hat{\mu}_{p, q}(-1) \partial_2 \hat{\tau}(\hat{\Pi}_p^{-1}(0, 0, \eta, \hat{\eta})).$$

Indeed, the value of this first term can be easily changed by a C^∞ -small localized perturbation. For example, we can keep the first jet along the $W_{X, \text{loc}}^u(p)$ the same, hence $\hat{\mu}_{p, q}(-1)$ will stay the same for $\hat{F}$. At the same time we can perturb the second jet at $\hat{\Pi}_p^{-1}(0, 0, \eta, \hat{\eta})$ changing the derivative of Jacobian at this point with respect to ξ , thus changing the value of $\partial_2 \hat{\tau}(\hat{\Pi}_p^{-1}(0, 0, \eta, \hat{\eta}))$ (indeed, what we are actually perturbing here is the Jacobian of the diffeomorphism $\bar{F}$, because the roof function is the logarithmic Jacobian + const). $\square$

Remark 3.13. Note that usage of coarse chart (and corresponding template and series $\hat{P}_p$) is really helpful for making the perturbation. Otherwise, if we use linearizing chart Φ_p , then perturbing the Jacobian would necessarily result in perturbation of the chart (because it needs to be linearizing for the perturbation), which makes the calculation very delicate.

Lemma 3.14. *Assume that the diffeomorphism $\hat{F}$ and its associated suspension flow satisfy $\hat{\mathcal{T}}_p^{ws}(\hat{q}) - \hat{P}_p(\hat{q}) \neq 0$. Then the same holds for any C^2 -small perturbation G of $\hat{F}$, and for all $(\eta', \hat{\eta}')$ sufficiently close to $(\eta, \hat{\eta})$, namely, the functions $\hat{\mathcal{T}}_{p^G}^{ws}$ and $\hat{P}_{p^G}$ for the associated suspension flow of G satisfy*

$$\hat{\mathcal{T}}_{p^G}^{ws}(\eta', \hat{\eta}') - \hat{P}_{p^G}(\eta', \hat{\eta}') \neq 0.$$

Proof. Recall that the coarser charts $\hat{\Phi}_{p^G}$ depend continuously on the flow in C^1 topology, and, accordingly, on the diffeomorphism in C^2 topology. The template $\hat{\mathcal{T}}_{p^G}^{ws}$ is just a particular coordinate (relative to a continuously varying chart) of the stable distribution which depends continuously on the flow in C^1 topology on the space of flows (hence, in C^2 topology on G); and $\hat{P}_{p^G}$ is given by an infinite series involving a first order derivative of the roof function (as in Subsection 3.2) which is given by the logarithm of the Jacobian of G , hence also varies continuously in C^2 topology with G . Thus, it is clear that the map $(G, (\eta', \hat{\eta}')) \mapsto \hat{\zeta}_{p^G}(\eta', \hat{\eta}') = \hat{\mathcal{T}}_{p^G}^{ws}(\eta', \hat{\eta}') - \hat{P}_{p^G}(\eta', \hat{\eta}')$ is continuous in C^2 topology on G . $\square$

3.4. Arriving at a contradiction via particular choices of a homoclinic point and parameters. In this section, we conclude the proof of the main results, i.e., Theorem A and Theorem B. We start with the proof of the deformation rigidity of typical perturbation of de la Llave examples, and then explain how it can be adapted to show the perturbative rigidity result.

According to the discussion in Section 2.1, all that is left to do is to prove Propositions 2.4-2.6, that is, verifying that given “generic” isospectral diffeomorphisms near de la Llave’s examples, the conjugacy between them preserves the strong stable/unstable manifolds at some periodic point. This follows immediately from Proposition 3.10 and the following propositions.

Proposition 3.15. *Consider an Anosov diffeomorphism G , a center-expanding periodic point $p^G = Y^T(p^G)$ for its suspension flow Y^t , and $(\eta, \hat{\eta}) \in (-1, 1)^2$, $r > 0$, as in Proposition 3.10, such that the functions $\hat{\mathcal{T}}_{p^G}^{ws}$ and $\hat{P}_{p^G}$ associated to Y^t satisfy*

$$\hat{\mathcal{T}}_{p^G}^{ws}(\eta', \hat{\eta}') - \hat{P}_{p^G}(\eta', \hat{\eta}') \neq 0, \quad \forall (\eta', \hat{\eta}') \in B((\eta, \hat{\eta}), r).$$

Then for any isospectral deformation $\{G_s\}_{s \in [0,1]}$ based at $G_0 = G$, we have $h_s(W_{G, \text{loc}}^{ss}(p^G)) = W_{G_s, \text{loc}}^{ss}(p_s^G)$, where h_s is the conjugacy map between G and G_s , and p_s^G is the continuation of p^G for G_s .

Proof. In the following, we slightly abuse notation and identify periodic points p_s^G for the diffeomorphisms G_s with the associated periodic points for the flows Y_s^t . For each $s \in [0, 1]$, we abbreviate as $\Sigma_s = \Sigma_{p_s^G}$ a transversal for the suspension flow Y_s^t of G_s containing the local manifolds $W_{Y_s^t, \text{loc}}^s(p_s^G)$, $W_{Y_s^t, \text{loc}}^u(p_s^G)$, we abbreviate as $\Phi_s = \Phi_{p_s^G}$ a chart in which the associated Poincaré map is linearized⁷ as in (2.2), we abbreviate as $\hat{\Phi}_s = \hat{\Phi}_{p_s^G}$ a coarse chart as in Subsection 3.2, and we denote by H_s the conjugacy between the suspension flows Y_0^t and Y_s^t ,

$$H_s \circ Y_0^t = Y_s^t \circ H_s, \quad \forall t \in \mathbb{R}.$$

Note that all these objects, except Φ_s , can be chosen to vary continuously with respect to the parameter $s \in [0, 1]$, cf. Remark 2.9. Moreover, by construction of the parametrizations Φ_s , and similarly for $\hat{\Phi}_s$, we have $H_s \circ \Phi_0((-1, 1)^2 \times \{(0, 0)\}) \subset \Phi_s((-1, 1)^2 \times \{(0, 0)\})$, because the local stable manifold $W_{Y_0^t, \text{loc}}^s(p^G) \subset \Sigma_0$ is mapped to the local stable manifold $W_{Y_s^t, \text{loc}}^s(p_s^G) \subset \Sigma_s$ by the conjugacy H_s .

Let us consider the set

$$\mathcal{G} := \{s \in [0, 1] : H_s(W_{Y_0^t, \text{loc}}^{ss}(p^G)) = W_{Y_s^t, \text{loc}}^{ss}(p_s^G)\}.$$

We want to show that the set $\mathcal{G}$ is equal to the whole interval $[0, 1]$. Given $\hat{\xi}_0 \in (-1, 1)$, we define Hölder continuous functions $[0, 1] \ni s \mapsto \hat{\xi}^\circ(\hat{\xi}_0, s)$ and $[0, 1] \ni s \mapsto \xi^\circ(\hat{\xi}_0, s)$ implicitly by

$$(\hat{\xi}^\circ(\hat{\xi}_0, s), \xi^\circ(\hat{\xi}_0, s), 0, 0) := \hat{\Phi}_s^{-1} \circ H_s \circ \hat{\Phi}_0(\hat{\xi}_0, 0, 0, 0), \quad \forall s \in [0, 1].$$

⁷Note that the linear map $\mathcal{L}$ does not depend on the parameter s due to the isospectrality condition.

We argue by contradiction and assume that $\mathcal{G} \neq [0, 1]$, i.e.,

$$\exists s_0 \in [0, 1] \quad \text{and} \quad \exists \hat{\xi}_0 \in (-1, 1) \quad \text{such that} \quad \xi^\circ(\hat{\xi}_0, s_0) \neq 0.$$

In the following, we fix such s_0 and $\hat{\xi}_0 \in (-1, 1)$ and abbreviate $\hat{\xi}^\circ(s) := \hat{\xi}^\circ(\hat{\xi}_0, s)$ and $\xi^\circ(s) := \xi^\circ(\hat{\xi}_0, s)$, for $s \in [0, 1]$. Without loss of generality, we can assume that $\xi^\circ(s_0) > 0$. By the continuity of strong stable manifolds, there exists a small strong stable segment $W_0 \subset W_{Y^t, \text{loc}}^{ss}(p^G)$ which contains the point $\hat{\Phi}_0(\hat{\xi}_0, 0, 0, 0)$ such that $H_{s_0}(W_0)$ is disjoint with the strong stable manifold $W_{Y_{s_0}^t, \text{loc}}^{ss}(p_{s_0}^G)$. See Figure 1. By continuity, there exist $\xi_0 < 0$, and $\varrho > 0$ such that for any $(\hat{\xi}', \xi') \in B((\hat{\xi}_0, \xi_0), \varrho)$, letting $(\hat{\xi}'(s), \xi'(s), 0, 0) := \hat{\Phi}_s^{-1} \circ H_s \circ \hat{\Phi}_0(\hat{\xi}', \xi', 0, 0)$, $s \in [0, 1]$, we have $\xi'(s) < 0$ for $s \approx 0$ and $\xi'(s) > 0$ for $s \approx s_0$.

By the density of homoclinic points, there exists a homoclinic point $q^G = \Phi_0(0, 0, \eta_\infty, \hat{\eta}_\infty) = \hat{\Phi}_0(0, 0, \eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in W_{Y^t, \text{loc}}^u(p^G)$, with $(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in B((\eta, \hat{\eta}), r)$, such that for some time $T' > 0$, the point $(q')^G := Y^{T'}(q^G)$ satisfies $(q')^G = \Phi_0(\hat{\xi}_\infty, \xi_\infty, 0, 0) = \hat{\Phi}_0(\hat{\xi}_\infty^\circ, \xi_\infty^\circ, 0, 0) \in W_{Y^t, \text{loc}}^s(p^G)$, with $(\hat{\xi}_\infty^\circ, \xi_\infty^\circ) \in B((\hat{\xi}_0, \xi_0), \varrho)$. For any $s \in [0, 1]$, let us consider the respective continuations q_s^G and $(q')_s^G$ of the points q^G and $(q')^G$:

$$\begin{aligned} q_s^G &= \Phi_s(0, 0, \eta_\infty(s), \hat{\eta}_\infty(s)) = \hat{\Phi}_s(0, 0, \eta_\infty^\circ(s), \hat{\eta}_\infty^\circ(s)) \in W_{Y_s^t, \text{loc}}^u(p_s^G), \\ (q')_s^G &= \Phi_s(\hat{\xi}_\infty(s), \xi_\infty(s), 0, 0) = \hat{\Phi}_s(\hat{\xi}_\infty^\circ(s), \xi_\infty^\circ(s), 0, 0) \in W_{Y_s^t, \text{loc}}^s(p_s^G). \end{aligned}$$

In particular, since $(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in B((\eta, \hat{\eta}), r)$, by Proposition 3.10, we have

$$\hat{\zeta}_{p^G}(q^G) = \hat{\mathcal{T}}_{p^G}^{ws}(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) - \hat{P}_{p^G}(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \neq 0,$$

and for $s \in (0, s_0)$, we have

$$\xi_\infty^\circ(s) < 0, \quad s \approx 0, \quad \xi_\infty^\circ(s) > 0, \quad s \approx s_0.$$

Since the map $s \mapsto \xi_\infty^\circ(s)$ is continuous (by the continuity of $s \mapsto \hat{\Phi}_s$), by the intermediate value theorem, we deduce that $\xi_\infty^\circ(s') = 0$, for some $s' \in (0, s_0)$. This means that at the parameter value s' the homoclinic point $(q')_{s'}^G$ lies precisely on the strong stable manifold of the point $p_{s'}^G$ as indicated on Figure 1 (recall item (2) from the definition of the coarse chart in Section 2.4).

Let us now explain how to reach a contradiction. On the one hand, by the (time-preserving) conjugacy between the flows Y_0^t and $Y_{s'}^t$, the periods of the shadowing periodic points p_n and their continuations for $Y_{s'}^t$ are equal to the same numbers T_n , which is independent of s (and the same applies to the excursion time T'). On the other hand, applying the asymptotic expansion given by Proposition 3.1 to the flow Y_0^t , we obtain

$$T_n - nT - T' = \omega_{p^G}(q^G, (q')^G) \mu^n + O(\theta^n),$$

with $\theta \in (0, \mu)$ and $\omega_{p^G}(q^G, (q')^G) := \xi_\infty(\mathcal{T}_{p^G}^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_{p^G}(\eta_\infty, \hat{\eta}_\infty)) \neq 0$, while for the flow $Y_{s'}^t$,

$$T_n - nT - T' = \omega_{p_{s'}^G}(q_{s'}^G, (q')_{s'}^G) \mu^n + O(\theta^n),$$

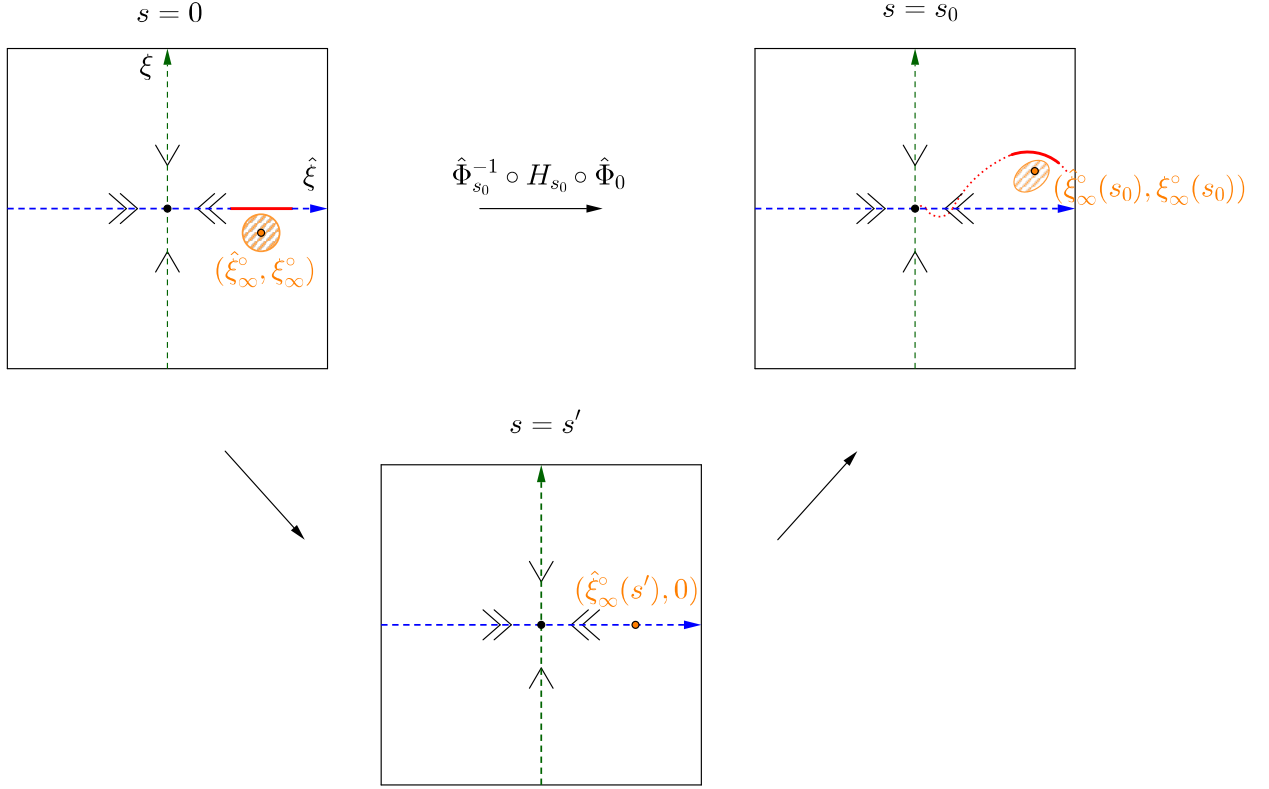
with $\omega_{p_{s'}^G}(q_{s'}^G, (q')_{s'}^G) := \xi_\infty(s')(\mathcal{T}_{p_{s'}^G}^{ws}(\eta_\infty(s'), \hat{\eta}_\infty(s')) - P_{p_{s'}^G}(\eta_\infty(s'), \hat{\eta}_\infty(s')))$. Comparing the two expressions above, we thus have

$$\omega_{p^G}(q^G, (q')^G) = \omega_{p_{s'}^G}(q_{s'}^G, (q')_{s'}^G) \neq 0.$$

By Corollary 3.9, we deduce that the associated quantity $\hat{\omega}_{p_{s'}^G}(q_{s'}^G, (q')_{s'}^G)$ in coarse charts also satisfies

$$\hat{\omega}_{p_{s'}^G}(q_{s'}^G, (q')_{s'}^G) = \xi_\infty^\circ(s') \zeta_{p_{s'}^G}(q_{s'}^G) \neq 0,$$

yielding a contradiction since s' was chosen such that $\xi_\infty^\circ(s') = 0$. We conclude that $\mathcal{G} = [0, 1]$, i.e., for any $s \in [0, 1]$, $H_s(W_{Y^t, \text{loc}}^{ss}(p^G)) = W_{Y_s^t, \text{loc}}^{ss}(p_s^G)$. Recall that for each $s \in [0, 1]$, the flow Y_s^t is a suspension flow over the Anosov diffeomorphism G_s ; more precisely, Y_s^t is the

FIGURE 1. Change of sign of the “center” coordinate $\xi_\infty^\circ(s)$ as s changes.

flow induced by the vertical flow in $\mathbb{T}^4 \times \mathbb{R}$ on the quotient manifold $M := (\mathbb{T}^4 \times \mathbb{R})/\hat{G}_s$, with $\hat{G}_s: (x, \alpha) \mapsto (G_s(x), \alpha - (K + \log J_s(x)))$. In particular, the local strong stable manifold $W_{Y_s^t, \text{loc}}^{ss}(p^G)$ lifts to a local graph in $\mathbb{T}^4 \times \mathbb{R}$ over the local strong manifold $W_{G_s, \text{loc}}^{ss}(p_s^G)$ for G_s . Moreover, the conjugacies H_s between the Anosov flows locally project to conjugacies h_s between the associated Anosov diffeomorphisms, from which we conclude that $h_s(W_{G_s, \text{loc}}^{ss}(p^G)) = W_{G_s, \text{loc}}^{ss}(p_s^G)$. $\square$

Let us now turn to the following perturbative version:

Proposition 3.16. *Consider an Anosov diffeomorphism G , a center-expanding periodic point $p^G = Y^T(p^G)$ for its suspension flow Y^t . Let $(\eta, \hat{\eta}) \in (-1, 1)^2$ and $r > 0$ be as in Proposition 3.10, such that the functions $\hat{\mathcal{T}}_{p^G}^{ws}$ and $\hat{P}_{p^G}$ associated to Y^t satisfy*

$$(3.9) \quad \hat{\mathcal{T}}_{p^G}^{ws}(\eta', \hat{\eta}') - \hat{P}_{p^G}(\eta', \hat{\eta}') \neq 0, \quad \forall (\eta', \hat{\eta}') \in B((\eta, \hat{\eta}), r).$$

Then there exists a C^2 open neighborhood $\mathcal{V}$ of G such that for any $G_1 \in \mathcal{V}$ which is isospectral to G , we have $h(W_{G_1, \text{loc}}^{ss}(p^G)) = W_{G_1, \text{loc}}^{ss}(p^{G_1})$, where h is the conjugacy map between G and G_1 , and p^{G_1} is the continuation of p^G for G_1 .

Proof. Without loss of generality, we can assume that the quantity in (3.9) is positive. By Lemma 3.14, we can thus fix a C^2 neighborhood $\mathcal{V}$ of G such that for any $G_1 \in \mathcal{V}$ conjugate to G , if we denote by Y_1^t the associated suspension flow, and by H the conjugacy between the suspension flows Y^t and Y_1^t , $H \circ Y^t = Y_1^t \circ H$, then we have

$$(3.10) \quad \hat{\mathcal{T}}_{p^G}^{ws}(H(\eta', \hat{\eta}')) - \hat{P}_{p^G}(H(\eta', \hat{\eta}')) > 0, \quad \forall (\eta', \hat{\eta}') \in B((\eta, \hat{\eta}), r).$$

Fix such a diffeomorphism $G_1 \in \mathcal{V}$. We take Y_1^t , H as above, and abbreviate $p = p^G$ and $p_1 = p^{G_1}$ in the following. We denote by Σ, Σ_1 transversals for the suspension flows Y^t, Y_1^t containing the local stable and unstable manifolds at p and p_1 , respectively. Let Φ, Φ_1 be

the associated linearizing charts, and let $\hat{\Phi}, \hat{\Phi}_1$ be the associated coarse charts. Assume by contradiction that

$$H(W_{Y^t, \text{loc}}^{ss}(p)) \neq W_{Y_1^t, \text{loc}}^{ss}(p_1).$$

Then, there exists $\hat{\xi}_0 \in (-1, 1)$ such that

$$(\hat{\xi}^\circ, \xi^\circ, 0, 0) := \hat{\Phi}_1^{-1} \circ H \circ \hat{\Phi}(\hat{\xi}_0, 0, 0, 0) \text{ has } \xi^\circ \neq 0.$$

Arguing as in the above proof of deformation rigidity, we can then find a homoclinic point $q^G = \Phi(0, 0, \eta_\infty, \hat{\eta}_\infty) = \hat{\Phi}(0, 0, \eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in W_{Y^t, \text{loc}}^u(p)$, with $(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in B((\eta, \hat{\eta}), r)$, such that for some time $T' > 0$, the point $(q')^G := Y^{T'}(q^G) \in W_{Y^t, \text{loc}}^s(p)$ satisfies

$$(q')^G = \Phi(\hat{\xi}_\infty, \xi_\infty, 0, 0) = \hat{\Phi}(\hat{\xi}_\infty^\circ, \xi_\infty^\circ, 0, 0) \in W_{Y^t, \text{loc}}^s(p), \text{ with } \xi_\infty^\circ < 0,$$

while for the points $q_1 := H(q^G) \in W_{Y_1^t, \text{loc}}^u(p_1)$ and $q'_1 := H((q')^G) = Y_1^{T'}(q_1) \in W_{Y_1^t, \text{loc}}^s(p_1)$, we have

$$q'_1 = \Phi_1(\hat{\xi}_\infty^1, \xi_\infty^1, 0, 0) = \hat{\Phi}_1(\hat{\xi}_\infty^{\circ, 1}, \xi_\infty^{\circ, 1}, 0, 0), \text{ with } \xi_\infty^{\circ, 1} > 0.$$

We then have the following expansions for the (common) periods T_n of the shadowing periodic points, for the flows Y^t and Y_1^t :

$$\begin{aligned} T_n - nT - T' &= \omega_p(q^G, (q')^G) \mu^n + O(\theta^n), \\ T_n - nT - T' &= \omega_{p_1}(q_1, q'_1) \mu^n + O(\theta^n), \end{aligned}$$

with $\theta \in (0, \mu)$, and

$$\begin{aligned} \omega_p(q^G, (q')^G) &:= \xi_\infty(\mathcal{T}_p^{ws}(\eta_\infty, \hat{\eta}_\infty) - P_p(\eta_\infty, \hat{\eta}_\infty)), \\ \omega_{p_1}(q_1, q'_1) &:= \xi_\infty^1(\mathcal{T}_{p_1}^{ws}(H(\eta_\infty, \hat{\eta}_\infty)) - P_{p_1}(H(\eta_\infty, \hat{\eta}_\infty))). \end{aligned}$$

In particular, we deduce that $\omega_p(q^G, (q')^G) = \omega_{p_1}(q_1, q'_1)$. Let $\hat{\omega}_p(q^G, (q')^G)$ and $\hat{\omega}_{p_1}(q_1, q'_1)$ be the corresponding quantities in coarse charts. Arguing as in the proof of Proposition 3.15, we deduce that $\hat{\omega}_p(q^G, (q')^G)$ and $\hat{\omega}_{p_1}(q_1, q'_1)$ have the same sign. But by the definition of the latter quantities, by (3.10) (recall that $(\eta_\infty^\circ, \hat{\eta}_\infty^\circ) \in B((\eta, \hat{\eta}), r)$), and since $\xi_\infty^\circ < 0 < \xi_\infty^{\circ, 1}$, we have $\hat{\omega}_p(q^G, (q')^G) < 0 < \hat{\omega}_{p_1}(q_1, q'_1)$, yielding a contradiction. $\square$

APPENDIX A. ON THE CONJUGACY BETWEEN 5-DIMENSIONAL ANOSOV FLOWS WITH A FINE SPLITTING INTO 1-DIMENSIONAL SUBBUNDLES

In this appendix, we introduce Theorem C, which extends the results previously established for the restricted class of suspension flows over perturbations of de la Llave's examples. While we omit the full proof of Theorem C, we outline the main ideas below. As we explain, some steps in the proof closely follow the arguments used for suspension flows, while the remaining points can be addressed using reasoning analogous to that in the proof of [GLH25, Theorem E].

Theorem C. Let $X^t: M \rightarrow M$ be a C^∞ transitive Anosov flow on a 5-manifold M which is not a constant roof suspension flow, with a fine dominated splitting into 1-dimensional subbundles:

$$TM = E_X^{ss} \oplus E_X^{ws} \oplus \mathbb{R}X \oplus E_X^{wu} \oplus E_X^{uu},$$

where $E_X^s := E_X^{ss} \oplus E_X^{ws}$ and $E_X^u := E_X^{wu} \oplus E_X^{uu}$ are uniformly contracted, respectively expanded, and $X := \frac{d}{dt}|_{t=0} X^t$ is the generating vector field of the flow. For $* \in \{ss, s, u, uu\}$, let W_X^* be the invariant foliation tangent to E_X^* . Assume that X^t is “center-dissipative”, i.e., $\log \det DX^t|_{E_X^{ws} \oplus E_X^{wu}}$ is not a coboundary. Then, at least one of the following points is true:

- (1) for each $x \in M$ the image of the stable subbundle E_X^s in the quotient bundle TM/E_X^u , restricted to $W_{X, \text{loc}}^u(x)$ has a C^∞ section;
- (2) for each $x \in M$ the image of the unstable subbundle E_X^u in the quotient bundle TM/E_X^s , restricted to $W_{X, \text{loc}}^s(x)$ has a C^∞ section;

- (3) there exists a C^1 -small neighborhood $\tilde{\mathcal{U}}$ of X such that if a C^∞ Anosov flow Y^t in the neighborhood $\tilde{\mathcal{U}}$ is C^0 conjugate to X^t , i.e., for some homeomorphism H , we have

$$(A.1) \quad H \circ X^t = Y^t \circ H, \quad \forall t \in \mathbb{R},$$

then:

- (a) X^t and Y^t have the same weak-stable/weak-unstable eigenvalues at periodic points, i.e., for any $p = X^T(p)$, and $* \in \{ws, wu\}$, we have $\det DX^T(p)|_{E_X^*} = \det DY^T(H(p))|_{E_Y^*}$;
- (b) H preserves strong foliations, i.e., for $* \in \{ss, uu\}$, $H(W_X^*(x)) = W_Y^*(H(x))$, $\forall x \in M$;
- (c) if, moreover, X^t and Y^t have the same strong-stable/strong-unstable eigenvalues at corresponding periodic points, then X^t and Y^t are C^{1+H} -conjugate.

Remark A.1. We stress that cases (1)-(2) of Theorem C are “exceptional”. Indeed, while the bundle TM/E_X^u is smooth along unstable leaves $W_{X,\text{loc}}^u(x)$, $x \in M$, the bundle E_X^s is in general only Hölder continuous along unstable leaves, hence we do not expect the existence of smooth sections as in case (1) (and similarly for case (2)). Note that while case (1) provided some additional partial smoothness of E_X^s , we do not claim that these local sections overlap in coherent way to give a global section. It is, however, natural to expect that with further arguments case (1) can be improved to a dichotomy: either image of E_X^s or the image of E_X^{ws} in TM/E_X^u is smooth along the unstable foliation.

Remark A.2. Let us also emphasize that the core of this paper focuses on showing items (3b)-(3c) within the framework of conjugate suspension Anosov flows over isospectral diffeomorphisms near de la Llave’s examples, under appropriate genericity conditions. In this context, the eigenvalue matching described in item (3a) is already ensured by the isospectrality of the underlying diffeomorphisms.

Below we give some details about the main steps of the proof of Theorem C. As in the statement of the theorem, we fix a 5-dimensional transitive Anosov flow X^t which is not a constant roof suspension flow.

In Subsection 2.3, given a periodic point $p = X^T(p)$ satisfying certain non-resonance conditions, we considered linearizing coordinates for the Poincaré map of a smooth transversal Σ_p and defined certain “templates” $\mathcal{T}_p^{ws}, \mathcal{T}_p^{ss}$ along its unstable manifold $W_{X,\text{loc}}^u(p)$. The issue is that these objects depend on the non-resonance conditions, hence are not uniform with respect to the periodic point p .

To handle this issue, similarly to the strategy in [GLH25], we replace linearizing coordinates with less nice normal forms, in order to keep uniformity with respect to the base point, following the construction of adapted charts in the work of Tsujii-Zhang [TZ23] (see also [GLH25, Proposition 3.2]). More precisely, we can construct a continuous family $\{\iota_x\}_{x \in M}$ of uniformly smooth charts such that the dynamics of X^t in these charts is partially normalized. We associate to these charts a family $\{\tilde{\Sigma}_x\}_{x \in M}$ of uniformly smooth sections transverse to the flow X^t , where for each $x \in M$, $\tilde{\Sigma}_x$ contains the local 2-dimensional stable and unstable manifolds $W_{X,\text{loc}}^s(x), W_{X,\text{loc}}^u(x)$; in particular, the jet of the hitting times of X^t along $W_{X,\text{loc}}^s(x), W_{X,\text{loc}}^u(x) \subset \tilde{\Sigma}_x$ is normalized to be a polynomial with uniformly bounded degree. In the same way as in Subsection 2.3, we can then define templates $\tilde{\mathcal{T}}_x^{ws}, \tilde{\mathcal{T}}_x^{ss}$ along $W_{X,\text{loc}}^u(x)$, and similarly, templates $\tilde{\mathcal{T}}_x^{wu}, \tilde{\mathcal{T}}_x^{uu}$ along $W_{X,\text{loc}}^s(x)$, where these objects are now defined relative to the sections $\tilde{\Sigma}_x$.

Then, following the proof of [GLH25, Proposition 4.2], given a center-expanding periodic point p , homoclinic points $q, q' = X^{T'}(q)$, and for the associated sequence $(p_n)_n$ of periodic points, we can derive asymptotic expansions of their periods T_n similar to those obtained in Proposition 3.1, but in terms of the templates $\tilde{\mathcal{T}}_x^{ws}$ corresponding to the uniform charts: the

formula in Proposition 3.1 would then become

$$T_n = nT + T' + \tilde{\xi}_\infty(\tilde{\mathcal{T}}_p^{ws}(\tilde{\eta}_\infty, \tilde{\eta}'_\infty) - \tilde{P}_p(\tilde{\eta}_\infty, \tilde{\eta}'_\infty))\mu^n + O(\theta^n),$$

where $\theta \in (0, \mu)$, where $\tilde{P}_p(\tilde{\eta}_\infty, \tilde{\eta}'_\infty)$ is a polynomial with uniformly bounded degree, and where $\tilde{\xi}_\infty, \tilde{\eta}_\infty, \tilde{\eta}'_\infty$ represent the coordinates of q, q' in $\tilde{\Sigma}_p$.

Instead of performing a perturbation as in Section 3.3 to ensure that the leading term $\tilde{\zeta}_p(q, q') := \tilde{\mathcal{T}}_p^{ws}(\tilde{\eta}_\infty, \tilde{\eta}'_\infty) - \tilde{P}_p(\tilde{\eta}_\infty, \tilde{\eta}'_\infty)$ in the above expansions is non-vanishing, we then argue as in [GLH25] and show a dichotomy:

- (1) either the term $\tilde{\zeta}_p(q, q')$ vanishes in a “robust way” (i.e., for a dense set of homoclinic points, by varying the points p, q, q'), in which case, we can conclude that the templates $\{\tilde{\mathcal{T}}_x^{ws}\}_{x \in M}$ are actually bounded polynomials; this corresponds to case (1) of Theorem C.

Let us outline the latter implication. Denote by $\bar{E}_X^s$ the image of E_X^s in the quotient bundle TM/E_X^u . Recall that given a point $x \in M$ and $y \in W_{X, \text{loc}}^u(x)$, the weak stable template $\tilde{\mathcal{T}}_x^{ws}(y)$ is defined so that

$$Z_x(y) := (0, 1, \tilde{\mathcal{T}}_x^{ws}(y), *, *) \in E_X^s(y).$$

Accordingly we have the local sections

$$(A.2) \quad \bar{Z}_x: W_{X, \text{loc}}^u(x) \rightarrow \bar{E}_X^s, \quad y \mapsto \bar{Z}_x(y) := (0, 1, \tilde{\mathcal{T}}_x^{ws}(y)) \in \bar{E}_X^s(y).$$

Such sections have the following properties:

- since $\{\tilde{\mathcal{T}}_x^{ws}\}_{x \in M}$ are bounded polynomials, for any $x \in M$, the section $\bar{Z}_x(\cdot)$ is uniformly smooth in $y \in W_{X, \text{loc}}^u(x)$;
 - $\text{span}(\bar{Z}_x(x)) = \bar{E}_X^{ws}(x)$, where $\bar{E}_X^{ws}$ is the image of E_X^{ws} in TM/E_X^u ;
 - the span of these sections is locally invariant, i.e., for $|t|$ sufficiently small, we have $D\bar{X}^t(y)(\bar{Z}_x(y)) \in \text{span}(\bar{Z}_{X^t(x)}(X^t(y)))$, where $D\bar{X}^t(y): \bar{E}_X^s(y) \rightarrow \bar{E}_X^s(X^t(y))$ is the map induced by the differential of the flow X^t between quotient spaces; indeed, in adapted charts, for small t , X^t simply becomes the vertical translation by t along the third coordinate (see [GLH25, Proposition 3.2]);
- (2) or for a set of “positive proportion” of center-expanding periodic points $p \in M$, we can find q, q' such that $\tilde{\zeta}_p(q, q') \neq 0$, in which case we can deduce several things as in case (3) of Theorem C:
 - (a) from period expansions, we can recover the weak-stable eigenvalue $\mu = \mu_p \in (0, 1)$ for this positive proportion set of periodic points p ; in particular, by the positive proportion Livshits theorem of [DMR24], if Y^t is another Anosov flow conjugate to X^t , with generating vector field sufficiently C^1 -close to the generating vector field of X^t , then we deduce that such eigenvalues have to match for the two flows as in item (3a) of Theorem C;
 - (b) if Y^t is another Anosov flow conjugate to X^t through a homeomorphism H , then, by the non-vanishing of the leading term $\tilde{\zeta}_p(q, q') \neq 0$ for suitable choices of p, q, q' , and since such property is robust (similarly to what was shown in Lemma 3.14), exactly as we did for suspension flows in the proof of Theorems A-B, if the generating vector field of Y^t is sufficiently C^1 -close to the generating vector field of X^t , then we can show that the conjugacy H sends the strong stable foliation W_X^{ss} of X^t to the strong foliation W_Y^{ss} of the conjugate flow Y^t as claimed in item (3b) of Theorem C;
 - (c) finally, from preservation of these foliations, and isospectrality conditions, we can deduce item (3c) of Theorem C in the same way as we did in the proof of Theorems A-B.

Of course, case (2) of Theorem C occurs in the symmetric way, when considering center-contracting periodic points instead of center-expanding periodic points. Similarly, information about strong unstable foliations can be deduced from the non-vanishing of the leading exponential term of associated period expansions.

Let us also stress that the case of “swapping SRB measures” which appears in the similar result [GLH25, Theorem E] does not occur here, namely in case (3a), due to the C^1 -closeness we assume on the vector fields X, Y (see, e.g., the proof of [GLH25, Theorem C] for more details).

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DEPARTMENT OF MATHEMATICS, THE OHIO STATE UNIVERSITY, COLUMBUS, OH 43210, USA.
Email address: gogolyev.1@osu.edu

ÉCOLE POLYTECHNIQUE, CMLS, ROUTE DE SACLAY, 91128 PALAISEAU CEDEX, FRANCE.
Email address: martin.leguil@polytechnique.edu